## Characterization of Cu<sup>2+</sup> Impurity Ion in K<sub>2</sub>SnCl<sub>6</sub>:Cu Single Crystal

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(Received 1 April 2018, Received in final form 16 May 2018, Accepted 17 May 2018)

K<sub>2</sub>SnCl<sub>6</sub> single crystals doped with Cu were grown by the evaporation of saturated solution. The electron magnetic resonance (EMR) of the Cu<sup>2+</sup> ion in the crystal has been investigated by using an X-band spectrometer. The rotation pattern in the crystallographic plane together with spin Hamiltonian parameters of Cu<sup>2+</sup> ion in K<sub>2</sub>SnCl<sub>6</sub>:Cu single crystal shows that the local site symmetry of the Cu<sup>2+</sup> ion is cubic at room temperature. The calculated spectroscopic splitting factor *g* of Cu<sup>2+</sup> impurity ion is 2.1908(±0.0007). The hyperfine constants of Cu<sup>2+</sup> impurity ion due to the <sup>63</sup>Cu and <sup>65</sup>Cu isotopes in K<sub>2</sub>SnCl<sub>6</sub>:Cu crystal are 0.83(±0.01)×10<sup>-4</sup> cm<sup>-1</sup> and 0.89(±0.01)×10<sup>-4</sup> cm<sup>-1</sup>, respectively. It turns out that the Cu<sup>2+</sup> ion substitutes for the K<sup>+</sup> ion without any nearby charge compensation.

Keywords : K<sub>2</sub>SnCl<sub>6</sub> single crystal, Cu<sup>2+</sup> EMR, local site symmetry, substitutional site

#### 1. Introduction

 $K_2SnCl_6$  is the prototype of the family of hexahalogenometallates of the general type  $A_2BX_6$  which crystallizes at room temperature in the cubic anti-fluorite structure with space group Fm3m [1]. The bonding between  $A^+$  and  $BX_6^{2-}$  octahedron is ionic and within the BX<sub>6</sub> octahedron covalent bonding dominates between B metals and X halogens. The  $K^+$  ion cages occupied by the  $BX_6^{2-}$ octahedra per unit cell in the hexahalogenometallates are alternately half occupied.

The electron configurations of the  $B^{4+}$  metal ions decide the magnetic structure of the compounds and also significantly influence the mechanism of the metal-ligand binding in the BX<sub>6</sub> octahedron. The investigations of the magnetic properties [2-4] and the bonding mechanisms [5, 6] of the single crystal were the essential theoretical and experimental parts of research. The results of the thermal and spectral measurements [7] and the study of the lattice dynamics have provided important information on the phase transitions mechanism and the lattice dynamics.

The mixed crystals of these hexahalogenometallates compounds have been investigated thoroughly especially with NQR and Raman scattering [8-10]. However, there are many difficulties in understanding the properties of the local structure in detail, particularly, in the temperature region around the structural phase transitions as well as in the intermediate phase that is in the narrow temperature region between the high- and low-temperature phases. Around the phase transition region, the dynamic and static properties of the lattice vary abruptly with temperature. Complex phenomena also come out through the overlap of many different effects of the lattice dynamics. The phase transition of  $K_2SnCl_6$  has been investigated in detail [11, 12].

Electron magnetic resonance (EMR) gives information on the local structure and site location of point defects as well as impurity ions. In these study, EMR spectra of  $Cu^{2+}$  ion in K<sub>2</sub>SnCl<sub>6</sub>:Cu single crystal obtained at an X-band are analyzed in terms of the effective spin Hamiltonian with the spin S = 1/2. The local site symmetry as well as EMR parameters of  $Cu^{2+}$  impurity ion is determined. The substitutional site of the  $Cu^{2+}$  ion is also discussed by considering the ionic size, charge state, and microscopic site symmetry.

# 2. Crystal Structure and Experimental Aspects

The K<sub>2</sub>SnCl<sub>6</sub> crystallize in the cubic with space group  $O_h^5$ /Fm3m consisting of non-connected SnCl<sub>6</sub>-octahedra on a face centered cubic (fcc) lattice and K atoms on a primitive cubic lattice with half the lattice constant at

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**Fig. 1.** Unit cell of cubic  $K_2SnCl_6$  single crystal. The black circles denotes the  $K^+$  cation and the large open circles the octahedral  $(SnCl_6)^{2-}$  anion.

room temperature [13]. The unit cell, i.e. primitive unit cell has four molecules with nine atoms per molecule (two K, one Sn, and six Cl atoms) as shown in Fig. 1. The basis consists of a  $(SnCl_6)^{2-}$  complex ion and two K<sup>+</sup> ions.

The crystal structure is an fcc with lattice constant a = 0.9986 nm at 280 °C. The bond lengths between atoms are Cl-Sn = 0.24106 nm, Cl-K = 0.35301 nm, Cl-Cl<sup>n</sup> = 0.34091 nm, Cl-Cl<sup>nn</sup> = 0.36491 nm at room temperature. The two phase transition temperatures at  $T_{c1} = 260.5$  K and at  $T_{c2} = 255$  K of K<sub>2</sub>SnCl<sub>6</sub> are found by X-ray neutron diffraction with single crystals and powders [14]. The crystal structure of high temperature phase above  $T_{c1}$  is cubic  $(O_h^5)$ . On cooling the crystal undergoes two successive phase transitions into an orthorhombic  $(D_{2h}^2)$  above  $T_{c2}$  and a monoclinic phase  $(C_{2h}^5)$  below  $T_{c2}$ .

The host crystal  $K_2SnCl_6$  doped with the Cu<sup>2+</sup> ion was grown by the vaporization of saturated solution. Potassium chloride (KCl) was melted in distilled water with 0.1 n HCl to yield saturated aqueous solution. The solution was then added to hydrogen hexachlorostannate (H<sub>2</sub>SnCl<sub>6</sub>) stoichiometrically to precipitate the powder of potassium hexachlorostannate (K<sub>2</sub>SnCl<sub>6</sub>).

The color of the powder obtained was snow-white. The powder was dissolved in distilled water and evaporated slowly to obtain single crystal. The  $K_2SnCl_6$  single crystal was grown up in 0.5-5 mm after a few days. The host crystal  $K_2SnCl_6$  is doped with  $Cu^{2+}$  (0.9 mol%).

We used an X-band JEOL (JES-RE2X) EMR spectrometer to get the  $Cu^{2+}$  EMR spectra in  $K_2SnCl_6$  single crystal. The rotation patterns of the  $Cu^{2+}$  resonance fields



**Fig. 2.** (Color online) Typical EMR spectra of  $Cu^{2+}$  ion in the K<sub>2</sub>SnCl<sub>6</sub>:Cu single crystal at room temperature. The resonance spectra consist of isolated  $Cu^{2+}$  EMR center with two sets of hyperfine structures due to the <sup>63</sup>Cu and <sup>65</sup>Cu isotopes and clustered  $Cu^{2+}$  EMR center.

in the crystallographic planes were obtained at microwave frequency of 9.446 GHz. Only one set of resonance line was measured for Cu<sup>2+</sup> in K<sub>2</sub>SnCl<sub>6</sub>:Cu single crystal. This EMR result agrees with the NQR frequency with a single line above  $T \ge T_{c1}$  [15, 16].

Figure 2 shows a typical EMR spectra of Cu<sup>2+</sup> impurity ion in K<sub>2</sub>SnCl<sub>6</sub>:Cu single crystal at room temperature. The fine structure of Cu2+ ion consists of only one resonance line because the effective spin S of the ion is 1/2. The fine structure is split into four lines each by the hyperfine interaction of  ${}^{63}$ Cu (I = 3/2, 69.17 % abundance) and  ${}^{65}$ Cu (I = 3/2, 30.83 % abundance) nuclei, respectively. The resonance spectra in Fig. 2 discussed later consist of isolated Cu<sup>2+</sup> EMR center with two sets of hyperfine structures (8 lines) due to the 63Cu and 65Cu isotopes as well as clustered Cu<sup>2+</sup> EMR center. We could not find any angular variation of electron magnetic resonance spectra for Cu<sup>2+</sup> ion on the crystallographic *aa*-plane. This indicates that the local site symmetry of the Cu<sup>2+</sup> ion K<sub>2</sub>SnCl<sub>6</sub>:Cu single crystal is cubic. From the rotation pattern, EMR parameters of the Cu2+ ion were determined.

#### 3. Effective Spin Hamiltonian and Analysis

The Cu<sup>2+</sup> ion has the electron configuration  $3d^9$  and is a *D*-state ion characterized by S = 1/2 and the ground multiplet  ${}^2D_{5/2}$ . The hyperfine structure (zero field splitting) results from the interaction between the nuclear spins of  ${}^{63}$ Cu and  ${}^{65}$ Cu isotopes and their electronic spin. The

experimental resonance fields can be analyzed with an effective spin Hamiltonian [17, 18]:

$$H = H_Z + H_{hyp} = \mu_B \boldsymbol{B} \cdot \boldsymbol{g} \cdot \boldsymbol{S} + \boldsymbol{S} \cdot \boldsymbol{A} \cdot \boldsymbol{S}, \qquad (1)$$

where the terms  $H_Z$  are the Zeeman interaction and  $H_{hyp}$  the hyperfine interaction.  $\mu_B$  is the Bohr magneton, **B** is the applied magnetic field, **g** is the spectroscopic splitting tensor, **S** is the electron spin, and **A** is the hyperfine tensor.

Magnetic resonance lines due to the hyperfine structure of  $Cu^{2+}$  in K<sub>2</sub>SnCl<sub>6</sub> recorded in the crystallographic planes are used for calculating the spectroscopic splitting parameter *g* and the hyperfine constant *A*. The calculated *g* value of the isolated  $Cu^{2+}$  impurity ions is 2.1908(±0.0007) as shown in Table 1, where spectroscopic splitting parameter *g* is isotropic. The impurities with no neighboring  $Cu^{2+}$  impurity ions within the range of interaction are denoted as isolated  $Cu^{2+}$  ions.

The well-resolved resonance spectra arise from the magnetically diluted  $Cu^{2+}$  impurity ions (isolated  $Cu^{2+}$  EMR center) in the host lattice generally. However, the eight-line spectra were not resolved well in our K<sub>2</sub>SnCl<sub>6</sub>:Cu crystal doped with 0.9 mol% Cu concentrations as shown in Fig. 2. At 0.9 mol% Cu concentrations, the intensity of hyperfine structure of isolated  $Cu^{2+}$  impurity ions is very weak and broad signal becomes dominant. These broad-line spectra in Fig. 2 can be explained by a new Cu center (clustered  $Cu^{2+}$  EMR center) arising from the higher Cu concentrations. The broad resonance line due to clustered Cu impurity ions is strongly influenced by their Cu magnetic moments between clustered Cu ions.

The observed hyperfine resonance lines of  $Cu^{2+}$  impurity ion in K<sub>2</sub>SnCl<sub>6</sub>:Cu crystal are very weak, but their presence is indicated clearly. Two sets of hyperfine lines due to the isotopes <sup>63</sup>Cu (I = 3/2) and <sup>65</sup>Cu (I = 3/2) are indicated by the four red-colored bars and by the four blue-colored bars in Fig. 2, respectively. The last term in the Hamiltonian eq. (1) represents the hyperfine interaction. Since the hyperfine interaction term is known to be much smaller than the Zeeman term in Eq. (1), it is sufficient to regard it as the first-order perturbation to the Zeeman term. The calculated hyperfine constants of Cu<sup>2+</sup> impurity ion due to the <sup>63</sup>Cu and <sup>65</sup>Cu isotopes in K<sub>2</sub>SnCl<sub>6</sub>:Cu crystal are  $A^{63} = 0.83 \ (\pm 0.01) \times 10^{-4} \ \text{cm}^{-1}$  and  $A^{65} = 0.89 \ (\pm 0.01) \times 10^{-4} \ \text{cm}^{-1}$  as shown in Table 1, respectively. The hyperfine constants are isotropic within the experimental accuracy both <sup>63</sup>Cu and <sup>65</sup>Cu nuclei in K<sub>2</sub>SnCl<sub>6</sub> crystal. The ratio  $A(^{65}Cu)/A(^{63}Cu)$  for Cu<sup>2+</sup> ion in K<sub>2</sub>SnCl<sub>6</sub> is 1.0723. This value (1.0723) is similar to the ratio of magnetic moments  $\mu(^{65}Cu)/\mu(^{63}Cu) = 1.0712$  [19], as it ought to be.

To identify the substitutional site location of the Cu<sup>2+</sup> ion in the K<sub>2</sub>SnCl<sub>6</sub> single crystal, several points are considered as follows. First, the ionic sizes of Cu<sup>2+</sup> and host ions in the K<sub>2</sub>SnCl<sub>6</sub> crystal are compared in Table 2. The ionic radius of the  $Cu^{2+}$  ion is about 0.071-0.087 nm depending on the coordination number. It is reasonable to assume that  $Cu^{2+}$  ion occupies the  $K^+$  ion instead of the  $Sn^{4+}$  because the ionic radius of the K<sup>+</sup> ion (0.151-0.152) nm) is larger than that of the  $Sn^{4+}$  ion (0.069-0.083 nm). On top of that, the charge state of the  $Cu^{2+}$  ion is closer to the  $K^+$  ion rather than that of the  $Sn^{4+}$  ion. Therefore, the  $Cu^{2+}$  ion can substitute for the K<sup>+</sup> ion with remotely charge compensation. Suppose that Cu<sup>2+</sup> ion resides at the Sn<sup>4+</sup> site, which is unfavorable from the ion size, local distortion at the Sn<sup>4+</sup> site would arise because the ionic radius of the  $Sn^{4+}$  is smaller than that of the  $Cu^{2+}$  ion. The cubic local site symmetry of the Cu<sup>2+</sup> ion obtained from our EMR analysis favors the assumption that the Cu<sup>2+</sup> ion occupies the K<sup>+</sup> ion at the cubic symmetry. Therefore, it is reasonable to conclude that the Cu<sup>2+</sup> ion replaces the K<sup>+</sup> ion without local charge compensation at the nearby surroundings.

#### 4. Results

 $K_2$ SnCl<sub>6</sub> single crystals doped with Cu<sup>2+</sup> ion were grown by the vaporization of saturated solution. The rotation

**Table 1.** The spectroscopic splitting parameter g and hyperfine constants A of the isolated  $Cu^{2+}$  impurity ion in a K<sub>2</sub>SnCl<sub>6</sub>:Cu single crystal. The unit of A is  $10^{-4}$  cm<sup>-1</sup>.

g-factor	$A(^{63}Cu)$	$A(^{65}Cu)$	$A(^{65}Cu) / A(^{63}Cu)$	$\mu(^{65}Cu) / \mu(^{63}Cu)$
2.1908(±0.0007)	0.83(±0.01)	0.89(±0.01)	1.0723	1.0712

Table 2.	The	ionic	radii	of th	ne Cu <sup>2</sup>	+ ion	and	the	host	cations	in	the	K <sub>2</sub> SnCl <sub>6</sub>	single	crystal.
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ions	Cu	1 <sup>2+</sup>	k	+	Sn <sup>4+</sup>		
Coord. Number	4	6	4	6	4	6	
radii (nm)	0.071	0.087	0.151	0.152	0.069	0.083	

patterns of Cu<sup>2+</sup> ion are observed on the crystallographic *aa*-plane at room temperature. The hyperfine constants *A* as well as the spectroscopic splitting factor *g* of the Cu<sup>2+</sup> EMR center have been determined with the effective spin Hamiltonian. The isotropic *g* and *A* values together with rotational patterns of Cu<sup>2+</sup> impurity ion show that the Cu<sup>2+</sup> ion reside at the cubic symmetry site. This EMR study enables us to conclude that the local site symmetry of the Cu<sup>2+</sup> center in K<sub>2</sub>SnCl<sub>6</sub> single crystal is cubic. It is proposed that the Cu<sup>2+</sup> ion in the crystal substitutes for the K<sup>+</sup> ion without local charge compensation at surroundings close by.

### Acknowledgement

We would like to appreciate the crystal growing and helpful discussions of Professor Y. M. Seo, D. Choi, and S. Song.

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