# Effect of Heat-treatment Temperature on the Formation of ε-Fe<sub>2</sub>O<sub>3</sub> Nanoparticles Encapsulated by SiO<sub>2</sub>

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 $\epsilon$ -Fe<sub>2</sub>O<sub>3</sub> has received attention with particular interest because of its large coercive field at room temperature, high-frequency millimeter-wave absorption, and the coupling of its magnetic and dielectric properties. This work investigated the effect of heat treatment on the formation of  $\epsilon$ -Fe<sub>2</sub>O<sub>3</sub>/SiO<sub>2</sub> composites fabricated using reverse-micelle and sol-gel methods. The heating process was performed at various temperatures to figure out the optimal conditions for acquisition of the  $\epsilon$ -Fe<sub>2</sub>O<sub>3</sub> phase, which exhibits the largest coercive field among the Fe oxides. The sample treated at 1,075 °C had the highest percentage of  $\epsilon$ -Fe<sub>2</sub>O<sub>3</sub> phase, with a coercivity (H<sub>C</sub>) of 21.57 kOe measured at room temperature that reached a maximum of 23.7 kOe at 230 K. The measurement of the magnetization-temperature (M-T) curve for this sample also reveals the characteristic magnetic transition associated with  $\epsilon$ -Fe<sub>2</sub>O<sub>3</sub> within the temperature range of 40-150 K. The crystal structure of  $\epsilon$ -Fe<sub>2</sub>O<sub>3</sub> was confirmed using X-ray powder diffraction. Transmission electron micrographs revealed a broad size distribution of iron oxide nanoparticles ranging from 12 to 22 nm. The findings indicate that  $\epsilon$ -Fe<sub>2</sub>O<sub>3</sub> is a promising candidate with high electromagnetic-wave absorption capacity that is appropriate for high-speed wireless communication applications.

Keywords :  $\varepsilon$ -Fe<sub>2</sub>O<sub>3</sub>, reverse-micelle and sol-gel method, heat-treatment, coercivity

# 1. Introduction

Magnetic iron oxide-based nanoparticles have received considerable attention for specific applications [1-3]. Iron oxides are used in technological applications, such as permanent magnets, magnetic recording, and information storage [4, 5], as well as biomedical applications, including magnetic fluids, magnetofection, cancer therapy, and drug delivery [6-8]. They are also crucial to theoretical studies that examine the quantum tunneling of magnetization and the impact of interparticle magnetic interactions on a nanoparticle system's magnetic regime [9, 10]. There are three naturally occurring iron oxides: Fe(II) oxide (FeO), Fe(III) oxide (Fe<sub>2</sub>O<sub>3</sub>), and Fe<sub>3</sub>O<sub>4</sub> (compound of Fe(II)oxide and Fe(III)-oxide).

Iron (III) oxide exists as four polymorphs: α-Fe<sub>2</sub>O<sub>3</sub>

(hematite),  $\beta$ -Fe<sub>2</sub>O<sub>3</sub>,  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub> (maghemite), and  $\epsilon$ -Fe<sub>2</sub>O<sub>3</sub>; each polymorph has a distinct structure and properties [4, 11, 12]. Among these four polymorphs,  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> and  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub>, which occur in both bulk and nanosized forms, are commonly found in nature and have been widely investigated.  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> is a red-brown solid with corundum structure that exhibits excellent properties, including low cost, abundant availability, wide light absorption, environment compatibility, and thermal stability [2, 13, 14]. Accordingly,  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> is appropriate for many applications in catalysis [15], biomedicine and biotechnology [16], gas sensors [17], and rechargeable lithium-ion batteries [18, 19].  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub> has a cubic spinel structure and exhibits ferromagnetic ordering with a net magnetic moment (2.5  $\mu_B$  per formula unit) and high Neel temperature (~950 K). Additionally, its chemical stability and low cost enables wide application in magnetic recording devices, electromagnetic absorbers, and biomedical engineering [20-22]. In contrast,  $\beta$ -Fe<sub>2</sub>O<sub>3</sub> and  $\epsilon$ -Fe<sub>2</sub>O<sub>3</sub> are rare phases with low natural abundances. It is challenging to synthesize them as a single phase because they only exist as nanosized

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structures and are thermally unstable; they readily transform into the  $\alpha$ -phase. Notably,  $\epsilon$ -Fe<sub>2</sub>O<sub>3</sub>, which appears as an intermediate phase between  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> and  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub>, has an extremely large magnetic coercive field of 20 kOe at room temperature [23]. Furthermore, it exhibits electromagnetic-wave absorption with a high resonance frequency, which can be used in information storage, magnetic-field tunable devices, and magnetic recording media [23-25]. The physical aspects of the  $\varepsilon$ -Fe<sub>2</sub>O<sub>3</sub> phase are poorly understood [26-29]. Many chemical and physical methods have been used to prepare highpercentage ɛ-Fe<sub>2</sub>O<sub>3</sub>-phase samples; for example, chemical vapor decomposition [30], flame spray pyrolysis [31], spray drying method [32], sol-gel method [33], pulsed laser deposition [34], sputtering [35], and thermal decomposition [36, 37].

Here, we report the synthesis of high percentage of  $\varepsilon$ -Fe<sub>2</sub>O<sub>3</sub>/SiO<sub>2</sub> nanoparticles by combining the reversemicelle and sol-gel methods. The reverse micelle process offers distinct advantages over alternative methodologies, notably enhanced precision in regulating particle dimensions, shapes, uniformity, and dispersal. In addition, the sol-gel process affords superior control over composition, primarily due to the silica matrix's role in inhibiting the agglomeration and phase transition of E-Fe<sub>2</sub>O<sub>3</sub> particles into a-Fe<sub>2</sub>O<sub>3</sub>. During the reverse micelle procedure, the precursor Fe(OH)<sub>3</sub> was uniformly produced and well dispersed. The hydrolysis occurs slowly and homogenously in the sol-gel procedure. This combination method created favorable conditions for the formation of ε-Fe<sub>2</sub>O<sub>3</sub> particles in the range of several tens nanometers in size. We investigated the formation and stability of the ε-Fe<sub>2</sub>O<sub>3</sub> phase in the temperature range between 700 °C and 1,300 °C. These measurements established a maximal coercivity (H<sub>C</sub>) of ~21.57 kOe for the sample heated at 1,075 °C, which was identified as the optimal treatment temperature to obtain  $\epsilon$ -Fe<sub>2</sub>O<sub>3</sub>. The temperature dependence of H<sub>C</sub> was also investigated.

## 2. Experimental

Figure 1 schematically illustrates the synthetic route to  $\varepsilon$ -Fe<sub>2</sub>O<sub>3</sub> nanoparticles. In the reverse-micelle step, two types of reverse-micelle systems were prepared: Rmicelle A and B. With stirring, the two reverse-micelle solutions were prepared from the 0.03/0.12/0.33/1 molar ratio of cetyltrimethylammonium bromide/1-butanol/noctane/H<sub>2</sub>O. To R-micelle A, 0.74 mmol of Fe(NO<sub>3</sub>)<sub>3</sub> was added to obtain a yellow solution (R-A); to R-micelle B, 30 mmol of NH<sub>3</sub> was dissolved to obtain the R-B solution. The R-B solution was added dropwise into the R-A solution while stirring for 30 min to obtain a brown solution. Then, 6.7 mmol of tetraethoxysilane was slowly injected into the brown solution. Stirring was maintained for 24 h at room temperature, producing a vellow-brown solution. The obtained precipitate was collected by centrifugation, washed several times with CHCl<sub>3</sub> and CH<sub>3</sub>OH, and dried. Then, the dry powder was heated at 700 °C, 975 °C, 1,000 °C, 1,025 °C, 1,050 °C, 1,075 °C, 1,150 °C, or 1,300 °C for 4 h to obtain a series of samples.

The crystallinity of each sample was confirmed using high-resolution X-ray diffraction (XRD; Bruker D8; Bruker Corp., Billerica, MA, USA). The magnetic properties of the samples were evaluated using a vibrating sample magnetometer (model 7304; Lake Shore Cryotronics, Inc., Westerville, OH, USA) and a Quantum Design physical property measurement system (Evercool II-9T;



Fig. 1. (Color online) Schematic illustration of the synthesis of the  $\epsilon$ -Fe<sub>2</sub>O<sub>3</sub> phase.



Fig. 2. (Color online) X-ray diffraction patterns of samples heated at different temperatures.

Quantum Design, Inc., USA). The morphologies of the samples were examined by transmission electron microscopy (TEM; JEM 2010; JEOL Ltd.).

## 3. Results and Discussion

Figure 2 presents the XRD patterns of the samples heated at different temperatures. When heating at a low temperature (700 °C), there was only a broad peak near 21.7°, which was attributed to amorphous SiO<sub>2</sub>. Iron oxide phases formed with increasing temperature. Up to 1,075 °C, the many peaks observed were matched to  $\varepsilon$ -Fe<sub>2</sub>O<sub>3</sub> signature XRD peaks. The remaining several peaks with extremely low intensity that corresponded to the  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> phase indicated that the sample was not composed of 100 %  $\varepsilon$ -Fe<sub>2</sub>O<sub>3</sub> phase. This result has been frequently reported in other studies because the E-Fe<sub>2</sub>O<sub>3</sub> phase is thermally unstable [29, 38, 39]. An increasing temperature caused the intensity of the  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> peaks to significantly increase, indicating that the portion of  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> was increasing. The crystallite size of the  $\varepsilon$ -Fe<sub>2</sub>O<sub>3</sub> nanoparticles was estimated from the full-width at halfmaximum (FWHM) of the peak centered at  $2\theta = 32.93^{\circ}$ according to the Debye-Scherrer's equation, as follows:

$$\mathbf{D} = \frac{k\lambda}{\beta\cos\theta}$$

where k is the particle geometry-dependent constant,  $\lambda$  is the X-ray wavelength (1.5406 Å),  $\beta$  is the FWHM of the peak, and  $\theta$  is the diffracted angle. Here, the particle size of 1,075 °C-sample was estimated to be 16.78 nm.

The sizes and shapes of the  $Fe_2O_3$  particles were additionally investigated by TEM for the sample heated at



**Fig. 3.** (a) low-magnification and (b) high-magnification transmission electron micrographs of  $Fe_2O_3$  nanoparticles heated at 1,075 °C.

1,075 °C, which had the largest percentage of the epsilon phase. Silica matrices were etched by stirring in an aqueous sodium hydroxide solution at 70 °C for 24 hours. The phase ratio of the sample heated at 1,075 °C obtained using the reference intensity ratio method was 96.2 %  $\varepsilon$ -Fe<sub>2</sub>O<sub>3</sub> and 3.8 %  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>. The TEM images (Fig. 3) revealed spherical particles and a highly uniform particle size distribution ranging from 12 to 22 nm, with a mean size of  $17 \pm 1$  nm. This size range is favorable for maintenance of the  $\varepsilon$ -Fe<sub>2</sub>O<sub>3</sub> phase because it converts to the  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> phase if the particle size exceeds ~30 nm [12, 31]. The particle size observed by TEM was in good agreement with the size determined by XRD.

One of the signature magnetic properties of the  $\varepsilon$ -Fe<sub>2</sub>O<sub>3</sub> phase is its large  $H_c$ , which can be used as a unique criterion to determine whether the  $\varepsilon$ -Fe<sub>2</sub>O<sub>3</sub> phase exists in composite nanoparticles. Field-swept vibrating sample magnetometry measurement at room temperature was performed to obtain magnetic hysteresis loops of the samples, which revealed the field dependence of magnetization curves for samples heated at different temperatures. The sharp decrease at ~3,000 Oe indicated the presence of a minor amount of α-Fe<sub>2</sub>O<sub>3</sub>, which was an impurity in the  $\varepsilon$ -Fe<sub>2</sub>O<sub>3</sub> phase. The H<sub>C</sub> of samples increased with increasing heating temperature and reached a maximum at 1,075 °C. However, the signal of the  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> phase in the hysteresis loop persisted. A further increase in temperature reduced the H<sub>C</sub>. This behavior is in good agreement with the XRD result, indicating that the evolution of  $\varepsilon$ -Fe<sub>2</sub>O<sub>3</sub> and  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> is a function of temperature and the coexistence of two phases within a particular temperature range. According to the measurement, the highest  $H_C$  of ~21.57 kOe was achieved for the sample heated at 1,075 °C. To clearly visualize the change in H<sub>C</sub>, its amplitude was plotted as a function of heating temperature (Fig. 4(b)).



**Fig. 4.** (Color online) (a) Normalized magnetization versus magnetic-field curve for various heating temperatures of samples at 300 K. (b) Coercivity as a function of heating temperature. Inset illustrates  $Fe_2O_3/SiO_2$  nanoparticles.

The hysteresis loops of the sample heated at 1,075 °C were measured at 20, 50, 80, 100, 120, 150, 180, 200, 230, 250, 270, and 300 K in a full magnetic-field scan from -50 to 50 kOe. These loops revealed the temperature dependence of the magnetic properties (Fig. 5). As noted above, this sample continued to include some  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>; thus, the loops were a combination signal of  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> and  $\epsilon$ -Fe<sub>2</sub>O<sub>3</sub> phases. However, below 150 K,  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> is a weak ferromagnet; beyond this temperature, it is superparamagnetic, indicating that the  $H_C$  of  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> shows negligible variation according to temperature. For that reason, the change in H<sub>C</sub> was mostly related to the contribution of the  $\varepsilon$ -Fe<sub>2</sub>O<sub>3</sub>. The temperature dependency of  $H_{\rm C}$  is clearly shown in Fig. 6(a). The coercive field slightly increased with decreasing temperature, reaching a maximum value of 23.7 kOe at 230 K. However, a sharp decrease in H<sub>C</sub> was observed between 170 K and 100 K, with the lowest H<sub>C</sub> of 1.3 kOe occurring at 100 K. Additionally, H<sub>c</sub> slightly increased upon further cooling to 20 K, reaching a value of 6.7 kOe. This behavior was attributed to increasing anisotropy, which is most likely



**Fig. 5.** Magnetic hysteresis loops of the sample heated at 1,075 °C measured at different temperatures.

related to the dominance of surface anisotropy at low temperatures [40]. The abrupt decrease in H<sub>C</sub> measured at 100 K also confirmed the presence of  $\varepsilon$ -Fe<sub>2</sub>O<sub>3</sub>; this phenomenon originated from the magnetic phase transition, which led to a decrease in H<sub>C</sub>. The magnetic moment as a function of temperature is presented in Fig. 6(b). The zero-field cooling curve under 50 kOe increased at 41 K, achieving its maximum value at ~149 K, which constitutes behavior characteristic of the  $\varepsilon$ -Fe<sub>2</sub>O<sub>3</sub> phase.

### 4. Conclusion

High-concentration (~96.2 %)  $\epsilon$ -Fe<sub>2</sub>O<sub>3</sub> nanoparticles were synthesized via combined reverse-micelle and solgel methods. The structure properties of 1,075 °C-sample has been studied by XRD and TEM. These measurements



**Fig. 6.** (Color online) (a) Coercivity as a function of measurement temperature for the sample heated at 1,075 °C. (b) Zero-field cooling curve of the same sample measured at 50 kOe.

revealed roughly spherical nanoparticles with a mean size of ~17 nm. A substantial coercivity (H<sub>C</sub>) of 21.57 kOe and 23.7 kOe was observed at temperatures of 300 K and 230 K, respectively, for the sample subjected to heating at 1,075 °C. Based on the temperature dependence of the magnetic moment, a magnetic transition typical of  $\epsilon$ -Fe<sub>2</sub>O<sub>3</sub> was observed in the range of 40-150 K.

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