

Tunable Slow Light with Large Bandwidth and Low-dispersion in Photonic Crystal Waveguide Infiltrated with Magnetic Fluids

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Two kinds of magnetic fluids with different volume fractions are symmetrically filled into the W0.9 photonic crystal waveguide structure. The 2D plane-wave expansion method is used to investigate the slow light properties numerically. The constant group index criterion is employed to evaluate the slow light performance. The wavelength bandwidth $\Delta\lambda$ centering at $\lambda_0 = 1550$ nm varies from 32.4 to 44.2 nm when the magnetic field factor $\alpha_{||}$ changes from 0 to 1. And the corresponding normalized delay bandwidth product can be tuned from 0.221 to 0.258. For comparison and optimization, two infiltration cases are investigated and the more advantageous infiltration scheme is found.

Keywords : photonic crystal waveguides, slow light, tunability, magnetic fluid

1. Introduction

Photonic crystal waveguide is an effective means to achieve slow light, which can realize slow light with large and flexible bandwidth and thus has attracted great attention [1]. There are two main methods to obtain the tunable slow light with photonic crystal waveguides. One is to tune the slow light by adjusting the structure of the photonic crystal waveguide, such as engineering the hole shape or size, changing the period or direction of the photonic crystal lattice [2-5], using the coupled cavity or center hole in defect structure [6, 7]. The other is to control the refractive index of liquid infiltrated into the holes of the photonic crystal structure, for instance, infiltrating the slotted photonic crystal waveguide with microfluidic technique [8].

The second method is superior to the first one because tuning the slow light can be realized by changing the refractive index of the infiltrated liquid easily and timely. Magnetic fluid is a good candidate liquid, whose refractive index can be adjusted by externally magnetic field and the volume fraction of magnetic nanoparticle [9, 10]. For tuning the slow light, magnetic fluid is filled in the holes or used as the background of the photonic crystal

structure placed in the magnetic field. Through changing the magnetic field, i.e. changing the refractive index of magnetic fluid, the slow light parameter of the structure can be set at the desired values. This kind of magnetically controllable optical component has been proved feasible in two-dimensional model [11] and photonic crystal slab configuration [12]. Recently, Jahanbakhshian *et al.* have studied the slow light properties in photonic crystal waveguide with asymmetric microfluidic infiltration [13]. Two kinds of materials with different refractive indices are infiltrated into the photonic crystal structure asymmetrically, which implies an alternative way to tune the slow light. In this work, the 3D photonic crystal slab waveguide filled with two kinds of magnetic fluids (different in volume fractions) is proposed and the corresponding slow light properties are investigated. The structure and infiltration parameters are optimized to obtain tunable slow light with large bandwidth and low dispersion.

2. Waveguide Design

The structure we proposed is silicon-based photonic crystal slab with triangular lattice. The lattice constant is a and the radius of hole is $r = 0.3a$. The slab is sandwiched between silica and air with thickness of $h = 283$ nm (same to Ref. 11). Then, removing a row of holes and shifting other holes inward to form W0.9 line-defect waveguide.

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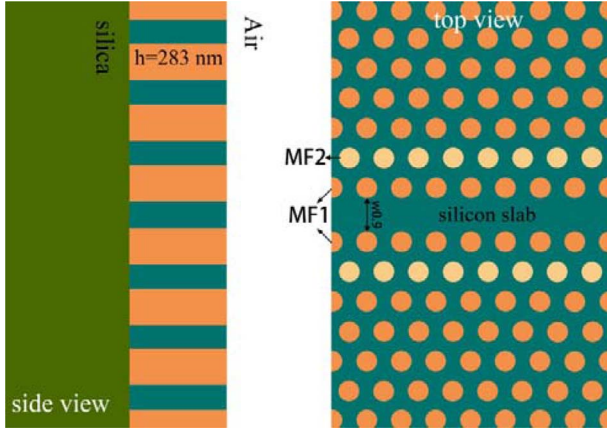


Fig. 1. (Color online) Schematic of the proposed photonic crystal waveguide structure. The lattice constant is a and $r = 0.3a$. Dark color of circle is MF₁ and light color is MF₂.

The details of waveguide structure are shown in Fig. 1. The magnetic fluid with volume fraction of ϕ_2 and refractive index of n_2 (referred as MF₂) is used to fill the second row of holes (see Fig. 1). The other rows of holes are filled with other magnetic fluid, whose volume fraction and refractive index are ϕ_1 and n_1 (referred as MF₁), respectively. Therefore, MF₁ and MF₂ are symmetrically filled into the photonic crystal structures on both sides of the line defect.

Through the Maxwell-Garnett approximation, the magnetic field dependent dielectric constant of magnetic fluid ϵ_{MF} can be expressed as [14-17]

$$\frac{\epsilon_{MF} - \epsilon_{liq}}{\alpha_{\parallel} \epsilon_{MF} + (3 - \alpha_{\parallel}) \epsilon_{liq}} = \phi \frac{\epsilon_{particle} - \epsilon_{liq}}{\epsilon_{particle} + 2 \epsilon_{liq}}, \quad (1)$$

where ϵ_{liq} is the dielectric constant of carrier liquid, $\epsilon_{particle}$ is the dielectric constant of magnetic nanoparticle, parameter ϕ is the volume fraction of magnetic nanoparticles in magnetic fluid, parameter α_{\parallel} means the degree of structural anisotropy caused by the magnetic field. α_{\parallel} changes in the range of 0-1 corresponding to the magnetic field varying from infinite to 0. $\epsilon_{particle} = 2.2^2$ and $\epsilon_{liq} = 1.448^2$ [12]. $\phi_1 = 0.5$ and $\phi_2 = 0.75$ are considered in current work. The magnetic permeability μ is taken as 1 at optical frequency. According to Eq. (1), the refractive indices of MF₁ and MF₂ are calculated out and listed in Table 1 and 2, respectively. The effective refractive index method based on 2D plane-wave expansion calculation is utilized to study the slow light performance [18, 19]. The refractive index of silica, silicon and air are 1.45, 3.48 and 1, respectively. Therefore, the equivalent slab effective index is obtained to be 3.0095.

Table 1. Refractive index of MF₁ for magnetic field factor α_{\parallel} varying from 0 to 1. $\epsilon_{liq} = 1.448^2$, $\epsilon_{particle} = 2.2^2$ and $\phi_1 = 0.5$.

Magnetic field factor α_{\parallel}	0	0.2	0.4	0.6	0.8	1
Refractive index n_1	1.7469	1.7555	1.7645	1.7741	1.7843	1.7952

Table 2. Refractive index of MF₂ for magnetic field factor α_{\parallel} varying from 0 to 1. $\epsilon_{liq} = 1.448^2$, $\epsilon_{particle} = 2.2^2$ and $\phi_2 = 0.75$.

Magnetic field factor α_{\parallel}	0	0.2	0.4	0.6	0.8	1
Refractive index n_2	1.8787	1.8968	1.9165	1.9381	1.9618	1.9879

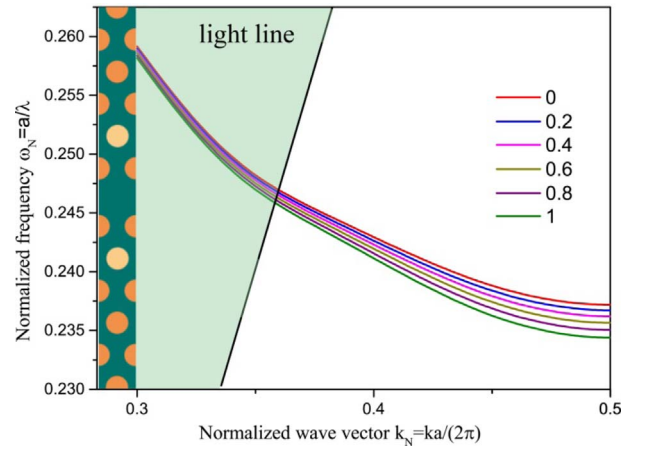


Fig. 2. (Color online) Dispersion curve of photonic crystal waveguide infiltrated with MF₁ and MF₂ symmetrically for magnetic field factor α_{\parallel} varying from 0 to 1. The supercell used for calculation is shown in the inset.

3. Results and Discussion

Using the data in Tables 1 and 2, the dispersion property of the guided mode in the photonic crystal waveguide is calculated and the result is shown in Fig. 2. Herein, the wave vector has been normalized as $k_N = ka/2\pi$, and the frequency has been normalized as $\omega_N = a/\lambda$. The light line of silica is also shown. Light can only propagate within the waveguide for the frequency under the light line. Fig. 2 indicates that the dispersion curve shifts to low frequency with the magnetic field factor α_{\parallel} increase from 0.

According to Fig. 2, the group index $n_g = c \frac{dk}{d\omega} = \frac{dk_N}{d\omega_N}$ is calculated out. The corresponding results are illustrated in Fig. 3. Fig. 3 shows that n_g is almost constant in the frequency range of $0.240 < \omega_N < 0.245$, which implies low-dispersion slow light.

According to $\beta_2 = \frac{d^2k}{d\omega^2} = \frac{1}{c} \frac{dn_g}{d\omega}$, the group velocity dispersion (GVD) β_2 is obtained from Fig. 3. The corresponding $\beta_2 \sim \omega_N$ curve is depicted in Fig. 4. Fig. 4 shows

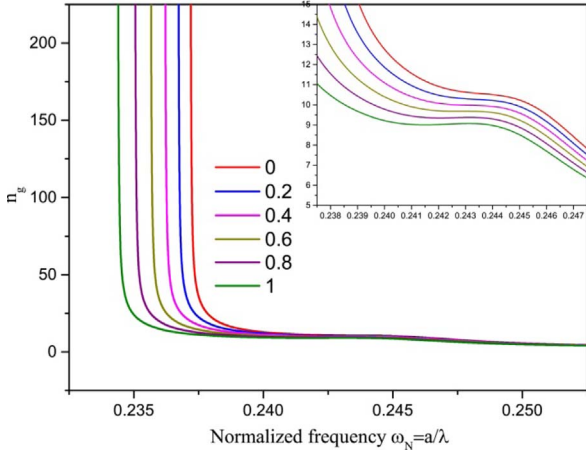


Fig. 3. (Color online) Group index of guided mode of the proposed structure at different magnetic fields.

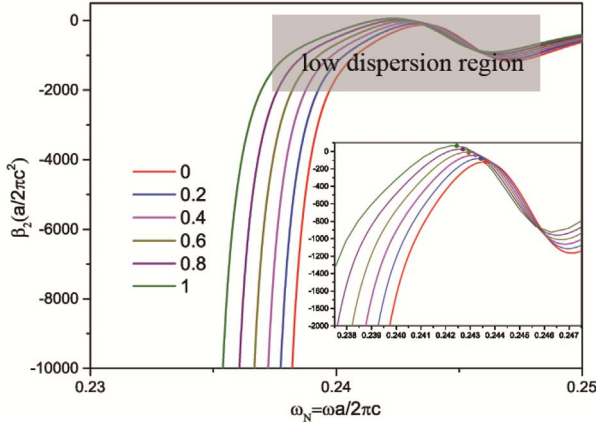


Fig. 4. (Color online) Group velocity dispersion of guided mode in photonic crystal waveguide filled with magnetic fluid for magnetic field factor α_{\parallel} varying from 0 to 1.

that there is a wide low-dispersion region with $-2 \times 10^3 a / (2\pi c^2) < \beta_2 < 100 a / (2\pi c^2)$. Comparing with the conventional low GVD definition, i.e. $\beta_2 < 10^6 a / (2\pi c^2)$ [20], the absolute value of β_2 in our low-dispersion region is nearly three orders of magnitude smaller. Therefore, our low-dispersion criterion is much stricter than that usually employed.

In order to further describe the slow light properties quantitatively and comparably, the normalized delay bandwidth product $NDBP = n_g \cdot \frac{\Delta\omega}{\omega_0} = n_g \cdot \frac{\Delta\omega_N}{\omega_{N_0}}$ is employed, where ω_{N_0} is center frequency, $\Delta\omega_N = \omega_{HN} - \omega_{LN}$. Herein, the constant group index criterion is used [21]. Therefore, the center frequency ω_{N_0} is selected at the peak of GVD curve (see the inset of Fig. 4), which ensures the slow light in the whole bandwidth $\Delta\omega_N$ has the lowest dispersion. Then, the constant group index n_g

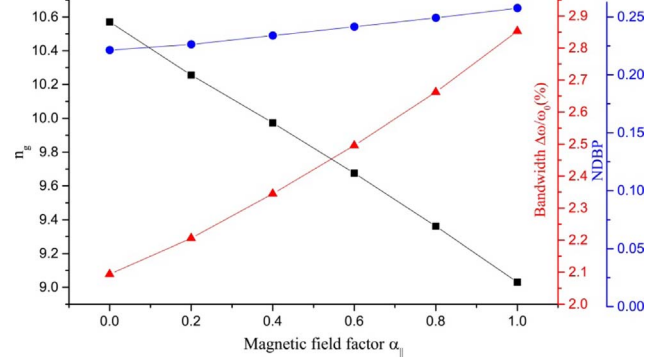


Fig. 5. (Color online) The relationship between magnetic field factor α_{\parallel} and group index, NDBP, and normalized bandwidth.

corresponding to ω_{N_0} can be obtained from Fig. 3. $n_{gH} = 1.1n_g$ and $n_{gL} = 0.9n_g$ corresponding to $\pm 10\%n_g$ are calculated. Thus, the frequencies ω_{HN} and ω_{LN} corresponding to n_{gH} and n_{gL} are obtained from Fig. 3. The relationship between magnetic field factor α_{\parallel} and group index n_g ,

NDBP and normalized bandwidth $\frac{\Delta\omega_N}{\omega_{N_0}}$ is plotted in Fig.

5. Fig. 5 indicates that the NDBP and bandwidth can be tuned with the magnetic field factor in the range of 0.221-0.258 and 2.09%-2.85%, respectively. So, if the center wavelength is selected at $\lambda_0 = 1550$ nm, the wavelength bandwidth $\Delta\lambda$ can be tuned from 32.4 to 44.2 nm when magnetic field factor α_{\parallel} changes from 0 to 1. Comparing with other optimized photonic crystal waveguides, our proposed photonic crystal waveguide has a much larger bandwidth [5, 22, 23].

In order to be comprehensive, the opposite case (viz. the second row of holes is filled with MF_1 and the other rows of holes are filled with MF_2) is also investigated. Numerical results indicate that the opposite case has lower slow light performance than that of the aforementioned one.

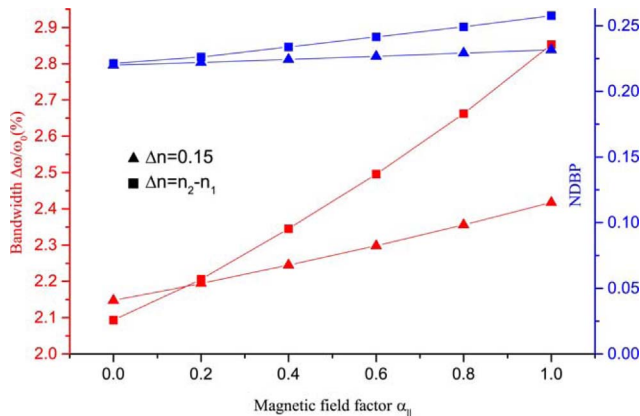
To further investigate the infiltration-dependent slow light properties, the magnetic-field-dependent and magnetic-field-independent $\Delta n = n_2 - n_1$ are considered, respectively. For the magnetic-field-dependent Δn case, the refractive indices listed in Table 1 and 2 are employed. According to Table 1 and 2, the relationship between Δn and magnetic field is listed in Table 3. For the magnetic-field-independent Δn case, $\Delta n = n_2 - n_1$ is set at 0.15. The refractive indices listed in Table 1 is employed as n_1 . So, $n_2 = n_1 + 0.15$ as listed in Table 4.

Table 3. Δn at different magnetic field factors.

Magnetic field factor α_{\parallel}	0	0.2	0.4	0.6	0.8	1
$\Delta n = n_2 - n_1$	0.1318	0.1413	0.1520	0.1640	0.1775	0.1927

Table 4. n_1 and n_2 at different magnetic field factors when $\Delta n = 0.15$.

Magnetic field factor α_{\parallel}	0	0.2	0.4	0.6	0.8	1
Refractive index n_1	1.7469	1.7555	1.7645	1.7741	1.7843	1.7952
Refractive index n_2	1.8969	1.9055	1.9145	1.9241	1.9343	1.9452

**Fig. 6.** (Color online) Bandwidth and NDBP for the magnetic-field-dependent and magnetic-field-independent Δn . The magnetic field factor α_{\parallel} varies from 0 to 1.

The slow light performances for the two cases are calculated and the corresponding results are shown in Fig. 6. Fig. 6 indicates, in terms of bandwidth and NDBP, the tunable range for the magnetic-field-dependent Δn case is larger than the magnetic-field-independent one. This implies that, besides changing refractive index of MF₁ and MF₂, changing Δn with magnetic field is more favorable for tuning the slow light of the proposed structure in a wide bandwidth and simultaneously with a large NDBP.

4. Conclusion

In summary, the W0.9 photonic crystal waveguide is designed and two types of magnetic fluids are proposed to infiltrate the photonic crystal structure in order to realize tunable slow light with large bandwidth and low dispersion. The second row of holes is filled with magnetic fluid MF₁ and other holes are filled with MF₂. The effective refractive index method based on 2D plane-wave expansion calculation is employed. The constant group index criterion is used to study the slow light properties. With the magnetic field factor α_{\parallel} varying from 0 to 1, the bandwidth can be tuned from 32.4 to 44.2 nm and NDBP can be tuned from 0.221 to 0.258. Besides, comparing with the magnetic-field-independent Δn case, tunable range of bandwidth and NDBP for the magnetic-field-dependent Δn case is superior. The obtained results imply

the advantages of the proposed structure, viz. tunable, low dispersion and high NDBP slow light.

Acknowledgments

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