

비접촉 동력전달용 자기장치의 성능예측을 위한 해석법

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- I Introduction
- II Analytical Approach - Coupling
- III Analytical Approach - Gear
- IV Conclusion

I

Introduction



Fig. Mechanical Couplings

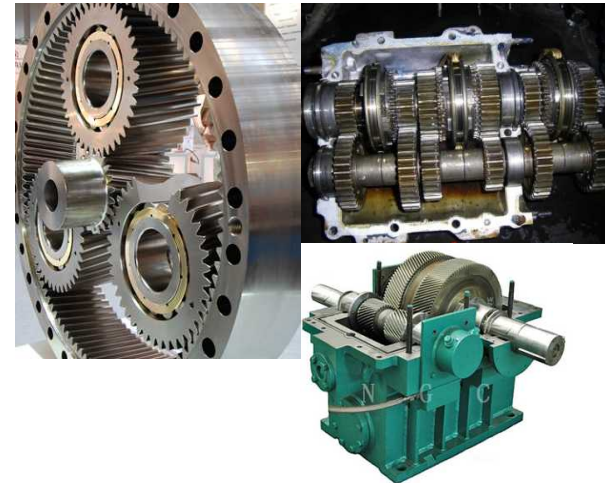


Fig. Mechanical Gears

- Couplings are mechanical devices for torque transmission in various industrial applications.
- Conventional couplings called mechanical coupling impart torque through mechanical contacts between separated parts.
- Maintenance is essential to prevent wearing down and destruction. Usually, mechanical couplings can be damaged in case of torque overload.

<http://www.alfabb.com/>
<http://www.ptreview.co.uk/>

I. Introduction



Fig. Broken mechanical couplings



Fig. Broken mechanical gears

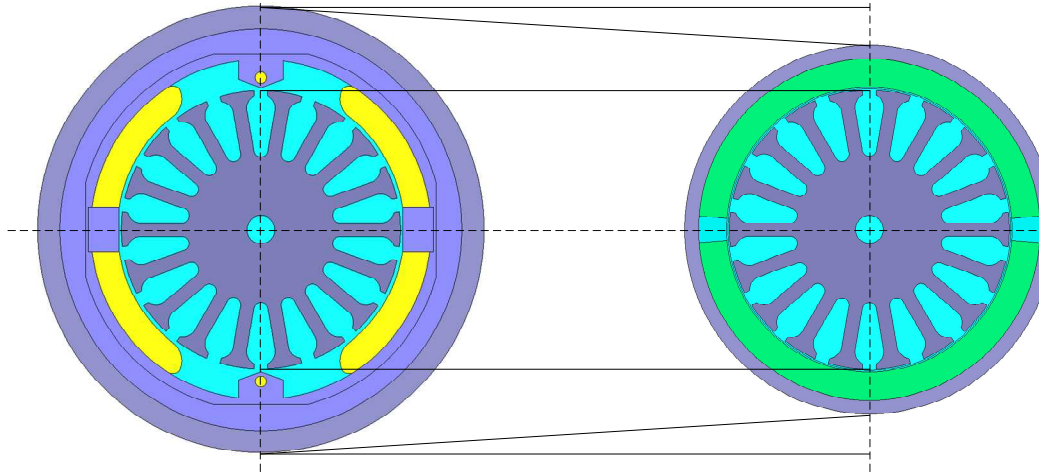
I. Introduction



Reference : Irving M. Gottlieb (1982) - "Electric Motors & Control Techniques"

DC

PMDC



Reference : Robert C. Perrine (1999) - "Design Handbook for PM Motors and Tachometers"

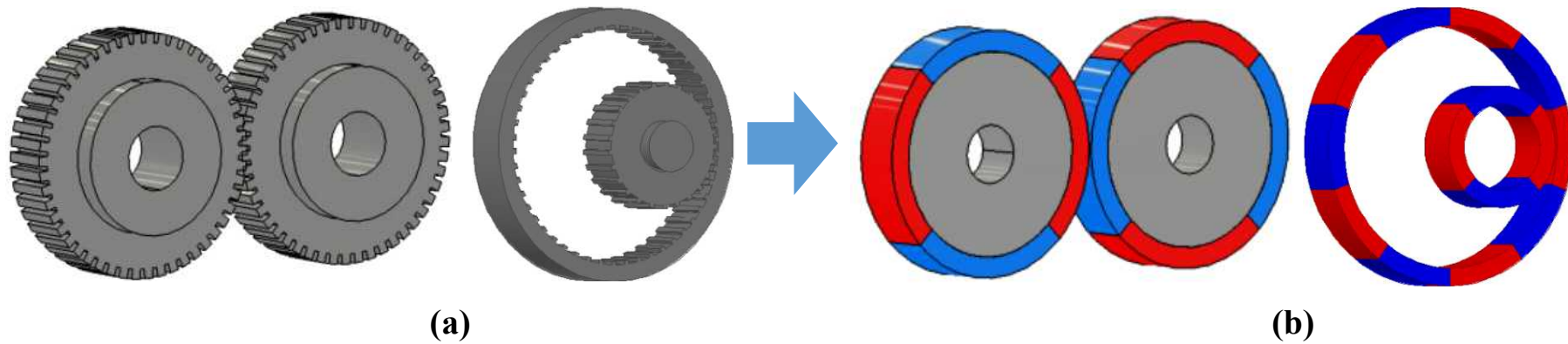


Fig. Structure of (a) Mechanical gears (b) Magnetic gears

- Advantages of magnetic gear

- ✓ Physical isolation
- ✓ Low maintenance
- ✓ Silent operation
- ✓ Inherent overload protection

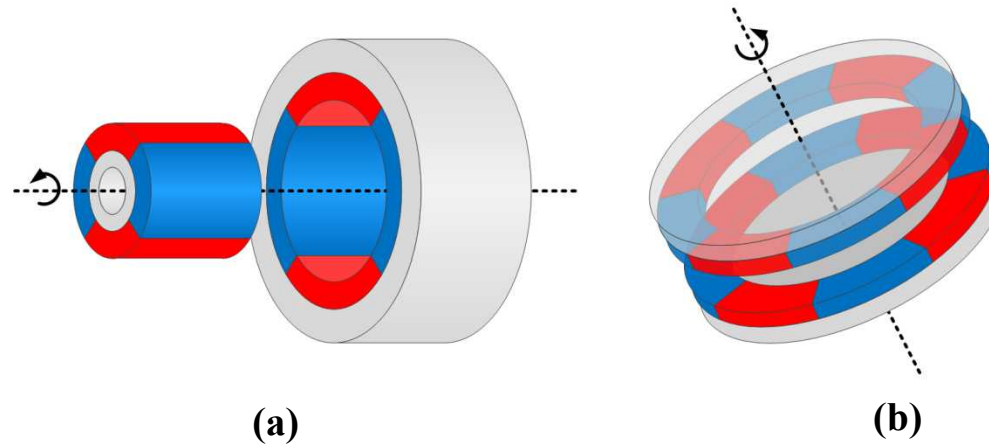


Fig. Structures of couplings: (a) RFPM (b) AFPM

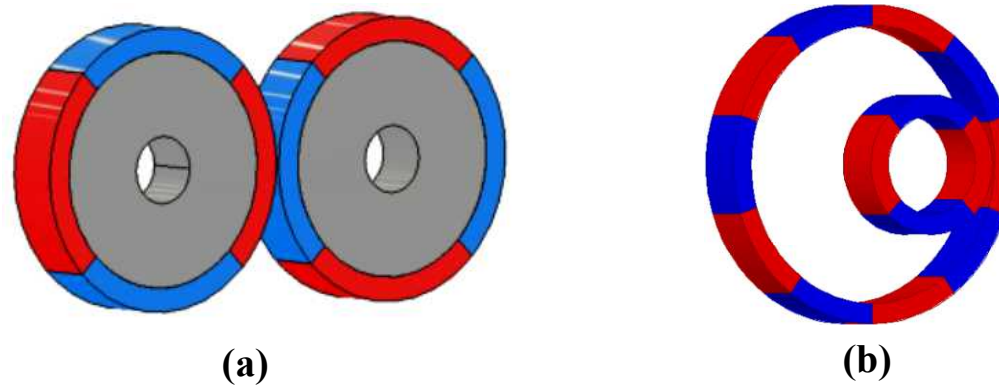


Fig. Structures of gears: (a) Spur (b) Cycloid

I. Introduction

Governing Equations

$$\mathbf{B} = \mu_0 (\mathbf{H} + \mathbf{M})$$

$$\nabla \times \mathbf{B} = \mu_0 \nabla \times \mathbf{H} + \nabla \times \mu_0 \mathbf{M}, \quad \xrightarrow{\nabla \times \mathbf{H} = \mathbf{J} = 0} \quad \nabla \times \mathbf{B} = \nabla \times \mu_0 \mathbf{M},$$

$$\mathbf{B} \equiv \nabla \times \mathbf{A}$$

$$\nabla \cdot \mathbf{A} = 0$$

$$\nabla^2 \mathbf{A} = -\mu_0 (\nabla \times \mathbf{M})$$

II Analytical Approach - Couplings

II. Analytical Approach – RFPM Coupling

A. Analytical Model

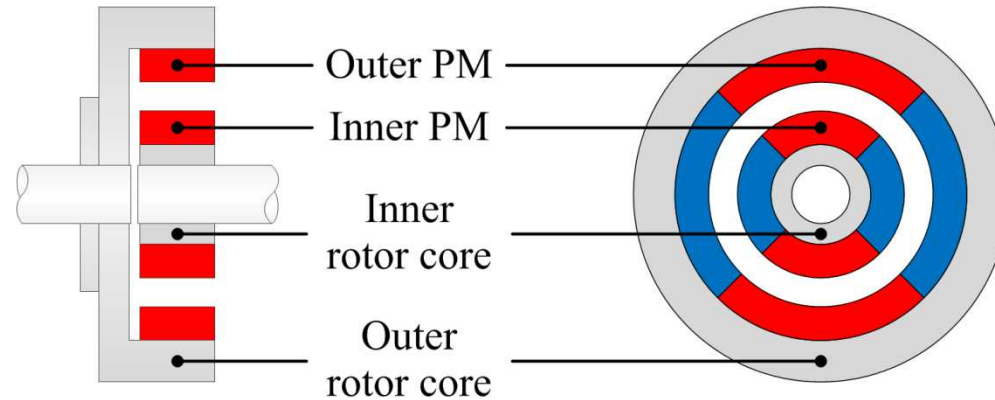


Fig. Schematic diagram of RFPM Coupling

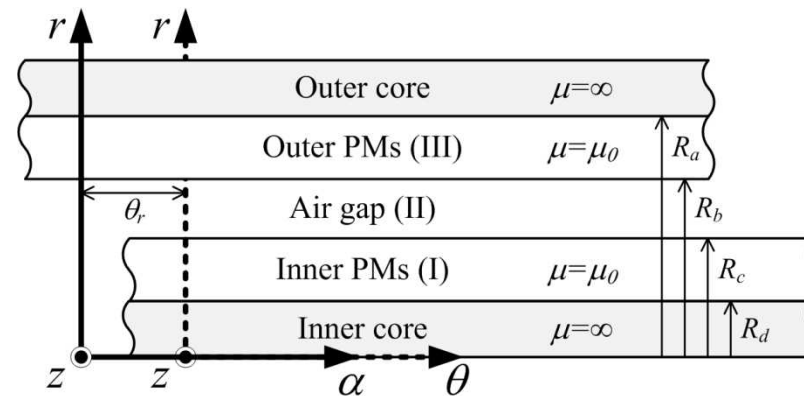


Fig. Analytical Model

II. Analytical Approach – RFPM Coupling

B. Magnetization

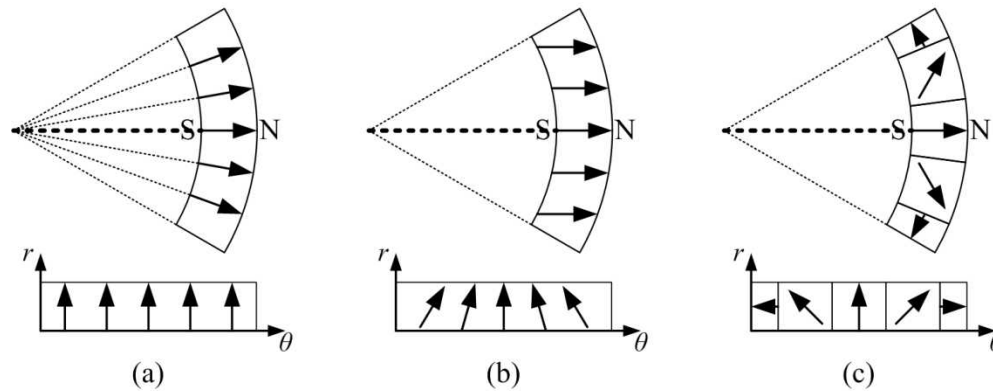


Fig. Possible Magnetization Patterns: (a) radial, (b) parallel, (c) Halbach

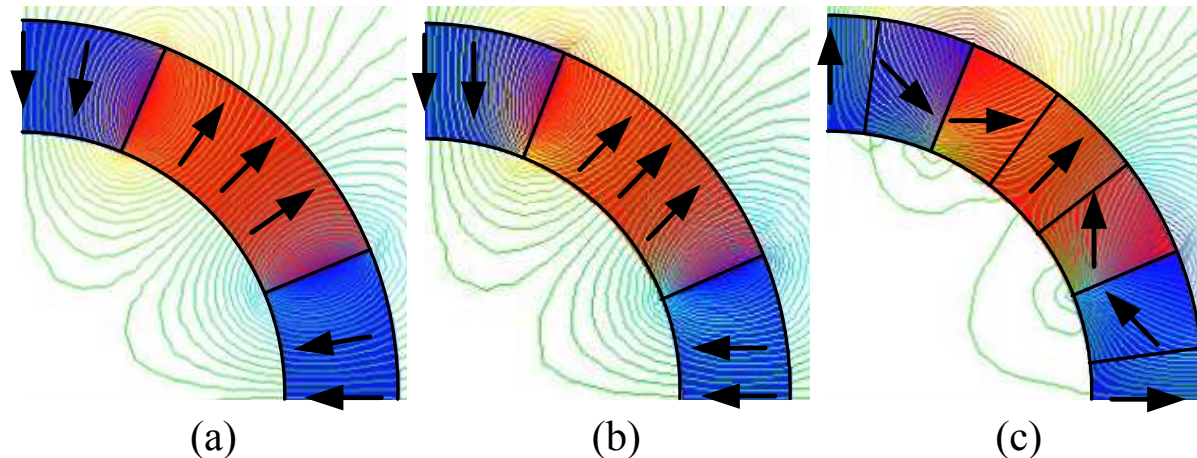


Fig. Magnetization Patterns : (a) radial magnetization (b) parallel magnetization (c) 3-segments Halbach magnetization

II. Analytical Approach – RFPM Coupling

B. Magnetization

1) Radial magnetization

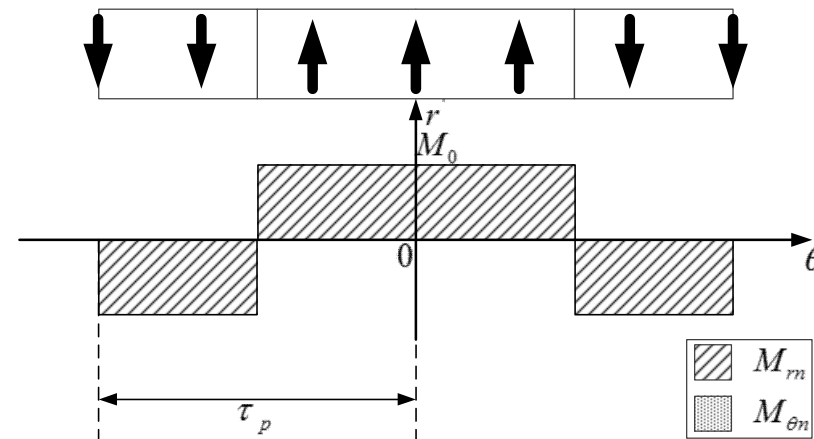


Fig. Radial magnetization model for Fourier series expansion

$$\vec{M} = \sum_{n=-\infty, \text{odd}}^{\infty} M_{rn} e^{-jn p_s \theta} \mathbf{i}_r$$

$$M_{rn} = \frac{jM_o}{n\pi} (e^{-jn\pi\theta_p/2p_s} - e^{jn\pi\theta_p/2p_s})$$

II. Analytical Approach – RFPM Coupling

B. Magnetization

2) Parallel magnetization

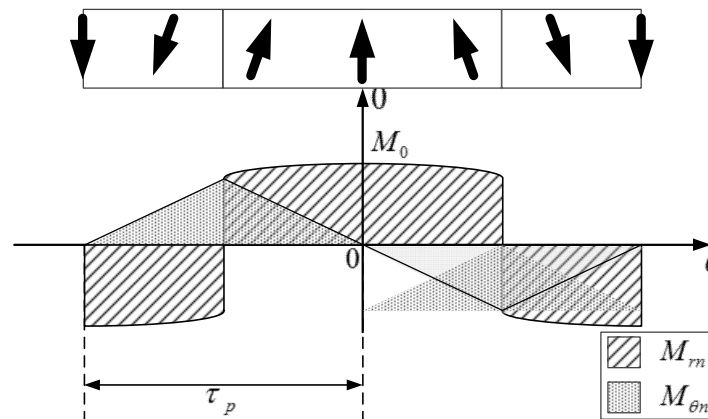


Fig. Parallel magnetization model for Fourier series expansion

$$\vec{M} = \sum_{n=-\infty, \text{odd}}^{\infty} (M_{rn} \mathbf{i}_r + M_{\theta n} \mathbf{i}_\theta) e^{-jnp_s \theta}$$

$$M_{\theta n} = \frac{M_0 p_s}{\pi \{1 - (np)^2\}} \begin{cases} e^{jnp_s \theta_p / 2} (\cos \frac{\theta_p}{2} - jnp \sin \frac{\theta_p}{2}) \\ -e^{-jnp_s \theta_p / 2} (\cos \frac{\theta_p}{2} + jnp \sin \frac{\theta_p}{2}) \end{cases}$$

$$M_{rn} = \frac{M_0 p_s}{\pi \{1 - (np)^2\}} \begin{cases} e^{jnp_s \theta_p / 2} (\sin \frac{\theta_p}{2} + jnp \cos \frac{\theta_p}{2}) \\ -e^{-jnp_s \theta_p / 2} (-\sin \frac{\theta_p}{2} + jnp \cos \frac{\theta_p}{2}) \end{cases}$$

II. Analytical Approach – RFPM Coupling

B. Magnetization

3) Halbach magnetization

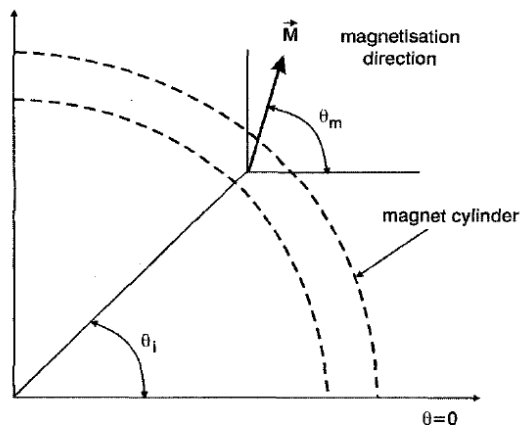


Fig.2 Angular relationship between θ_m and θ_i

$$\theta_m = (1 \pm p) \theta_i,$$

where θ_m and θ_i represent the direction of each individual magnet segment and the angle between criterion of angle ($\theta=0$) and the center of the i th magnet segment, respectively.

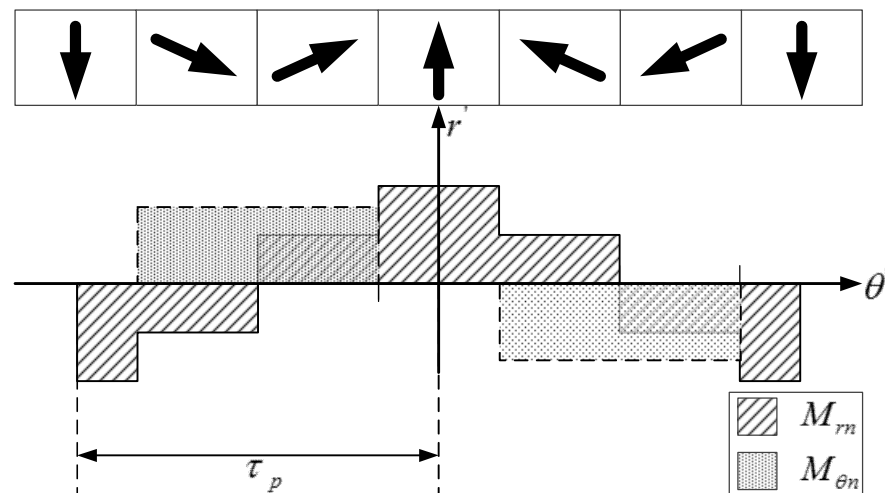


Fig. 3-segments Halbach magnetization model for a Fourier series expansion

$$\mathbf{M} = \sum_{n=-\infty, \text{odd}}^{\infty} (M_{rn} \cdot \mathbf{i}_r + M_{\theta n} \cdot \mathbf{i}_\theta) \cdot e^{-jnp_s \theta}$$

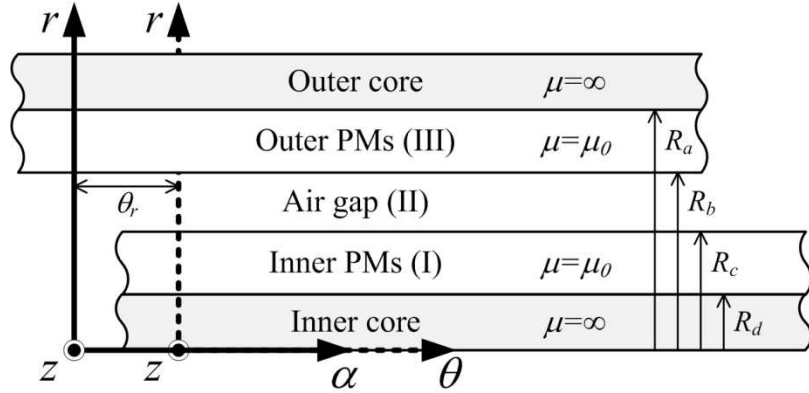
$$M_{rn} = \frac{M_0}{2\beta_p} \left\{ e^{\frac{5\beta_p}{6}} - e^{-\frac{5\beta_p}{6}} + e^{\frac{\beta_p}{6}} - e^{-\frac{\beta_p}{6}} \right\} + \frac{M_0}{\beta_p} \left\{ e^{-\beta_p} - e^{\beta_p} + e^{7\beta_p} + e^{5\beta_p} \right\}$$

$$M_{\theta n} = \frac{\sqrt{3}M_0}{4\beta_p} \left\{ e^{\frac{5\beta_p}{6}} + e^{-\frac{5\beta_p}{6}} - e^{\frac{\beta_p}{6}} - e^{-\frac{\beta_p}{6}} \right\}$$

Z. Q. Zhu and D. Howe, "Halbach permanent magnet machines and applications: a review," *IEE Proc.-Electr. Power Appl.*, vol. 148, pp. 299-308, Jul. 2001.

II. Analytical Approach – RFPM Coupling

B. Governing Equations



$$\frac{\partial^2}{\partial r^2} A_n^{I,III}(r) + \frac{1}{r} \frac{\partial}{\partial r} A_n^{I,III}(r) - \frac{(np)^2}{r^2} A_n^{I,III}(r) = -\frac{\mu_0 np}{r} M_n^{in,out}$$

$$\frac{\partial^2}{\partial r^2} A_n^{II}(r) + \frac{1}{r} \frac{\partial}{\partial r} A_n^{II}(r) - \frac{(np)^2}{r^2} A_n^{II}(r) = 0$$

$$\mathbf{A}^I = \sum_{n=-\infty, odd}^{\infty} \left\{ C_n^I r^{np} + D_n^I r^{-np} + \mu_0 r \frac{np}{(np)^2 - 1} M_n^{in} \right\} e^{-jnp\theta} \mathbf{i}_z$$

$$\mathbf{A}^{II} = \sum_{n=-\infty, odd}^{\infty} \left\{ C_n^{II} r^{np} + D_n^{II} r^{-np} \right\} e^{-jnp\alpha} \mathbf{i}_z$$

$$\mathbf{A}^{III} = \sum_{n=-\infty, odd}^{\infty} \left\{ C_n^{III} r^{np} + D_n^{III} r^{-np} + \mu_0 r \frac{np}{(np)^2 - 1} M_n^{out} \right\} e^{-jnp\alpha} \mathbf{i}_z$$

$$\nabla^2 \mathbf{A}^{I,III} = -\mu_0 (\nabla \times \mathbf{M}_{RF}^{in,out})$$

$$\nabla^2 \mathbf{A}^{II} = 0$$

$$\mathbf{A} = \sum_{n=-\infty, odd}^{\infty} A_n(r) e^{-jnp(\alpha \text{ or } \theta)} \mathbf{i}_z$$

$$\nabla \times \mathbf{M}_{RF}^{in,out} = \frac{1}{r} \begin{vmatrix} \mathbf{i}_r & r\mathbf{i}_\theta & \mathbf{i}_z \\ \frac{\partial}{\partial r} & \frac{\partial}{\partial \theta} & \frac{\partial}{\partial z} \\ \sum_{n=-\infty, odd}^{\infty} M_{rn}^{in,out} e^{-jnp(\alpha \text{ or } \theta)} & \sum_{n=-\infty, odd}^{\infty} r M_{\theta n}^{in,out} e^{-jnp(\alpha \text{ or } \theta)} & 0 \end{vmatrix}$$

$$= \sum_{n=-\infty, odd}^{\infty} \frac{np}{r} \left(\frac{M_{\theta n}^{in,out}}{np} + j M_{rn}^{in,out} \right) e^{-jnp(\alpha \text{ or } \theta)} \mathbf{i}_z$$

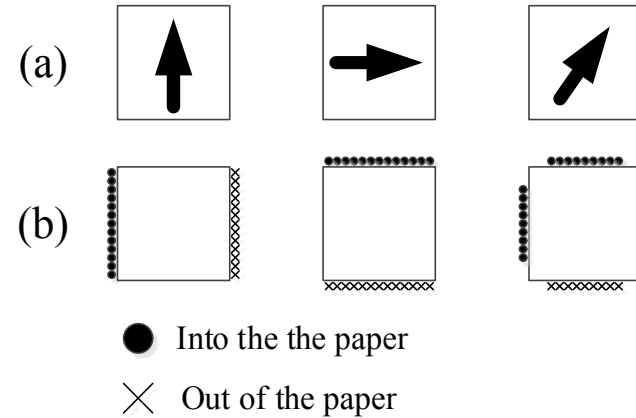
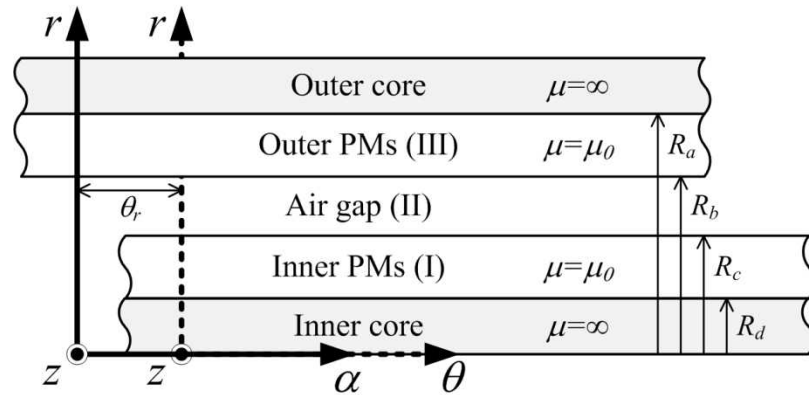
$$\mathbf{B} = \nabla \times \mathbf{A} = \frac{1}{r} \begin{vmatrix} \mathbf{i}_r & r\mathbf{i}_\theta & \mathbf{i}_z \\ \frac{\partial}{\partial r} & \frac{\partial}{\partial \theta} & \frac{\partial}{\partial z} \\ 0 & 0 & \sum_{n=-\infty, odd}^{\infty} A_n(r) e^{-jnp(\alpha \text{ or } \theta)} \end{vmatrix}$$

$$\mathbf{B}_\theta = - \sum_{n=-\infty, odd}^{\infty} \frac{\partial}{\partial r} A_n(r) e^{-jnp(\alpha \text{ or } \theta)} \mathbf{i}_\theta$$

$$\mathbf{B}_r = - \sum_{n=-\infty, odd}^{\infty} \frac{jnp}{r} A_n(r) e^{-jnp(\alpha \text{ or } \theta)} \mathbf{i}_r$$

II. Analytical Approach – RFPM Coupling

C. Boundary Conditions



$$H_{\theta n}^{III}(R_b, \alpha) - H_{\theta n}^{II}(R_b, \alpha) = M_{\theta n}^{out}$$

$$H_{\theta n}^I(R_c, \theta) - H_{\theta n}^{II}(R_c, \alpha) = M_{\theta n}^{in}$$

$$\mathbf{B}_r^{II}(R_b, \alpha) = \mathbf{B}_r^{III}(R_b, \alpha)$$

$$\mathbf{B}_r^I(R_c, \theta) = \mathbf{B}_r^{II}(R_c, \alpha)$$

$$\mathbf{B}_\theta^I(R_d, \theta) = 0$$

$$\mathbf{B}_\theta^{III}(R_a, \alpha) = 0$$

$$\begin{aligned} \mathbf{F} &= -S_{RF} \mu_0 \left\langle \mathbf{H}_r^{II}(R_b, \alpha, z) \mathbf{H}_\theta^{II}(R_b, \alpha, z) \right\rangle_\alpha \\ &= -\frac{S_{RF}}{\mu_0} \left[\mathbf{B}_r^{II}(R_b, \alpha, z) \cdot \left\{ \mathbf{B}_\theta^{II}(R_b, \alpha, z) \right\}^* \right] \end{aligned}$$

II. Analytical Approach – RFPM Coupling

D. Results

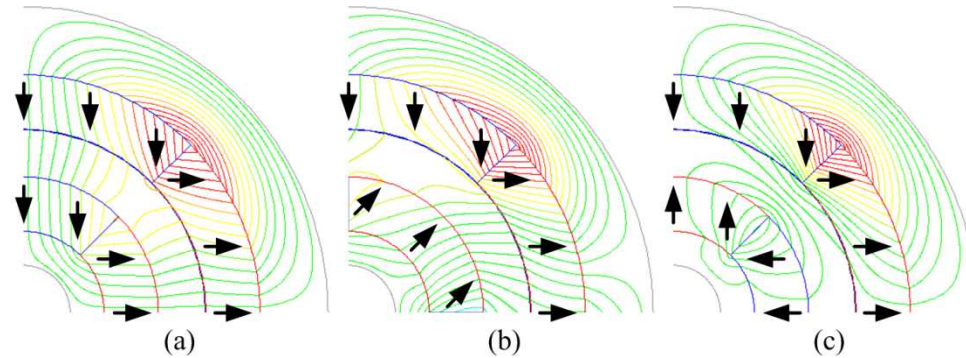


Fig. Magnetic flux line distribution of parallel RFPMC according to relative angular shift: (a) $\theta_r = 0^\circ$, (b) $\theta_r = 45^\circ$, and (c) $\theta_r = 90^\circ$.

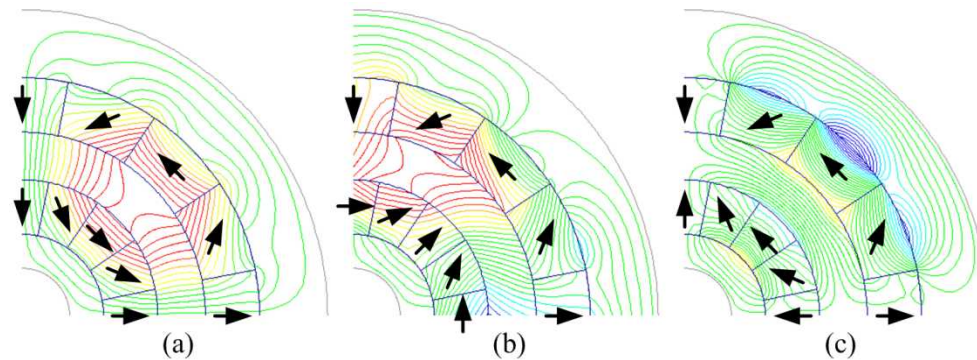


Fig. Magnetic flux line distribution of Halbach RFPMC according to relative angular shift: (a) $\theta_r = 0^\circ$, (b) $\theta_r = 45^\circ$, and (c) $\theta_r = 90^\circ$.

II. Analytical Approach – RFPM Coupling

D. Results

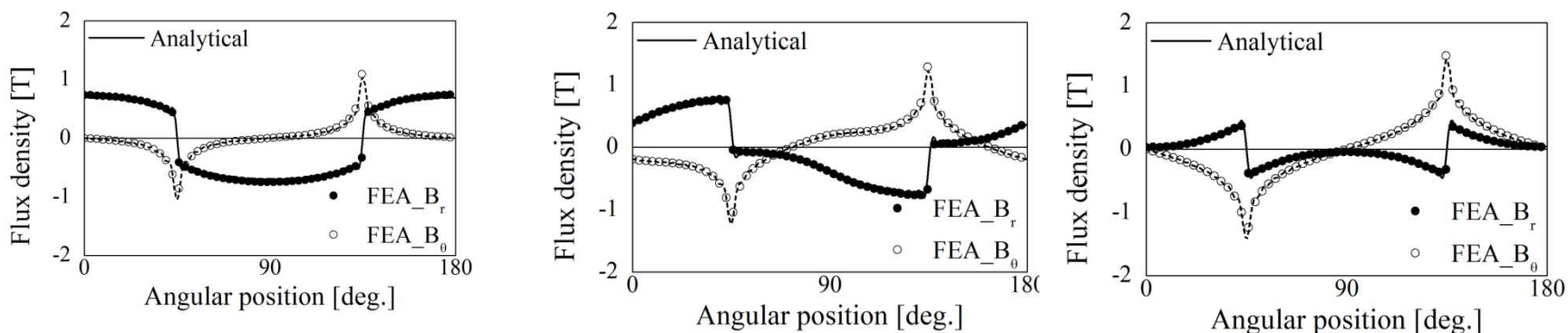


Fig. Comparison of analytical results with FE results for air gap flux density of parallel RFPMC when θ_r is 0, 45, and 90°.

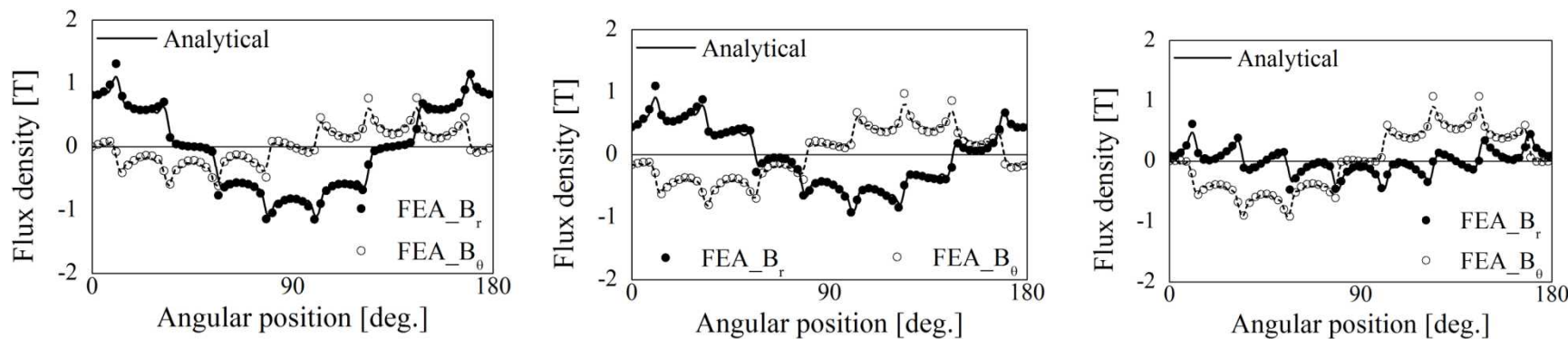


Fig. Comparison of analytical results with FE results for air gap flux density of Halbach RFMC when θ_r is 0°, 45, and 90.

II. Analytical Approach – RFPM Coupling

D. Results

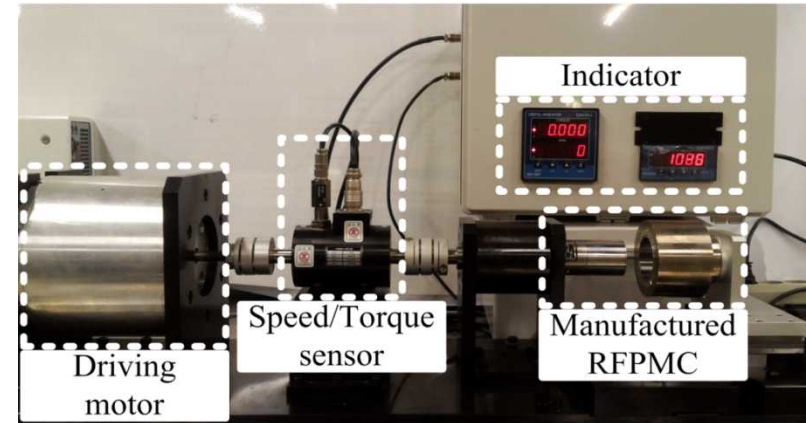
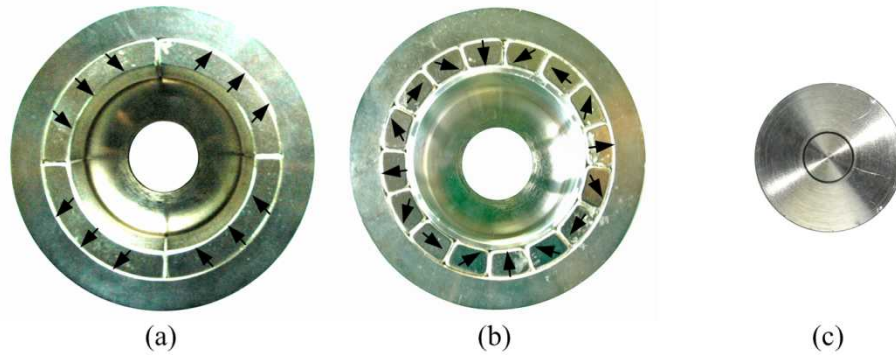
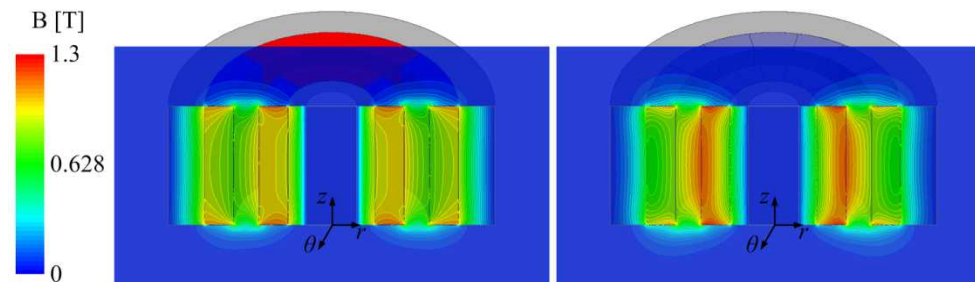


Fig. Actual manufactured RFPMCs; (a) parallel magnetized outer rotor, (b) Halbach magnetized outer rotor and (c) inner rotor.

		Analytical results	3D FEA results	Measured results
Parallel magnetization	Torque	16.49 [Nm]	14.71 [Nm]	15.05 [Nm]
	Error	9.6 %	2.3 %	
Halbach magnetization	Torque	19.06 [Nm]	16.72 [Nm]	17.08 [Nm]
	Error	11.6 %	2.1 %	



II. Analytical Approach – AFPM Coupling

A. Analytical Model

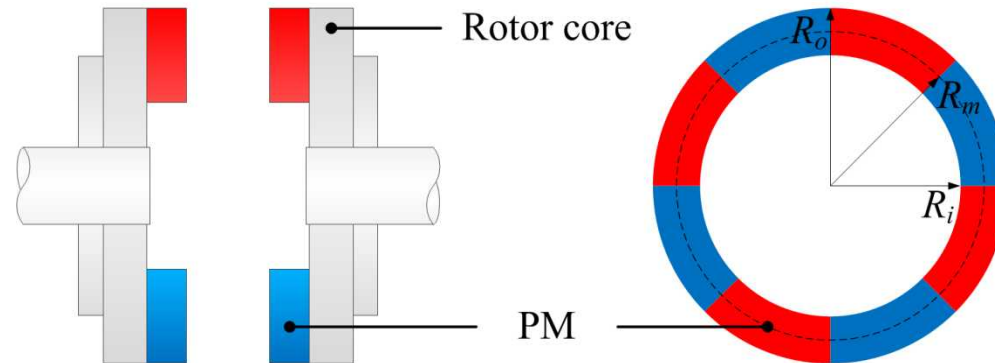


Fig. Schematic diagram of RFPM Coupling

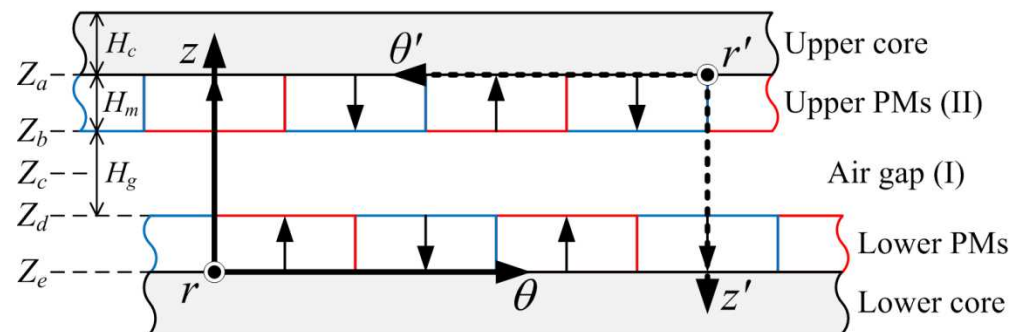


Fig. Analytical Model

II. Analytical Approach – AFPM Coupling

B. Magnetization

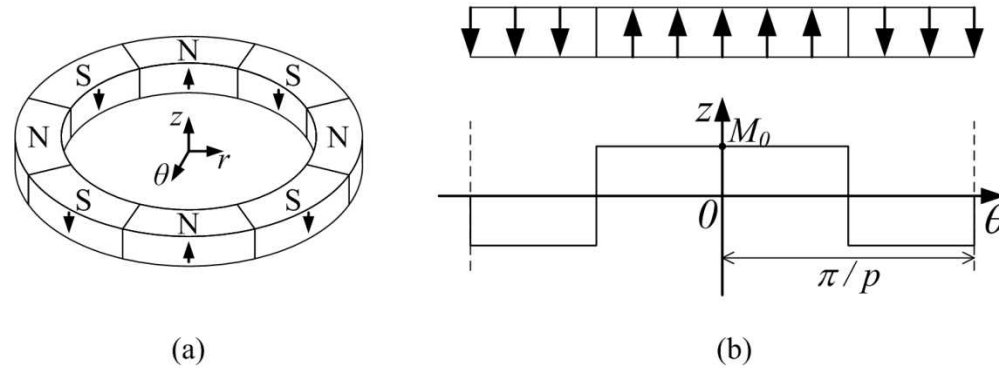


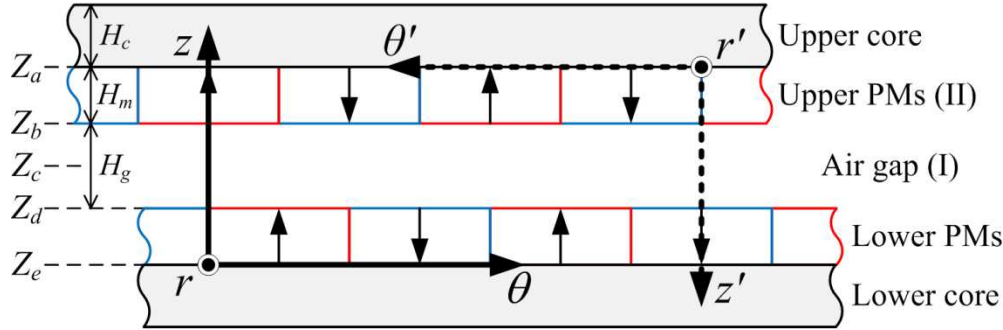
Fig. (a) Axial direction magnetization of AFPMC (b) axial magnetization model for Fourier series expansion.

$$\mathbf{M}_{AF} = \sum_{n=-\infty, \text{odd}}^{\infty} M_{zn} e^{-jnp\theta} \mathbf{i}_z$$

$$M_{zn} = \frac{M_0}{jn\pi} \left(e^{jnp\frac{\theta_w}{2}} - e^{-jnp\frac{\theta_w}{2}} \right)$$

II. Analytical Approach – AFPM Coupling

B. Governing Equations



$$\nabla^2 \mathbf{A}^{\text{II}} = -\mu_0 (\nabla \times \mathbf{M}_{\text{AF}})$$

$$\nabla^2 \mathbf{A}^{\text{I}} = 0$$

$$\mathbf{A} = \sum_{n=-\infty, \text{odd}}^{\infty} A_n(z) e^{-jnp\theta} \mathbf{i}_r$$

$$\mathbf{A}^{\text{I}} = \sum_{n=-\infty, \text{odd}}^{\infty} \left(E_n^{\text{I}} e^{\frac{np}{r}z} + F_n^{\text{I}} e^{-\frac{np}{r}z} \right) e^{-jnp\theta} \mathbf{i}_r$$

$$\mathbf{A}^{\text{II}} = \sum_{n=-\infty, \text{odd}}^{\infty} \left(E_n^{\text{II}} e^{\frac{np}{r}z} + F_n^{\text{II}} e^{-\frac{np}{r}z} - \frac{j\mu_0 r}{np} M_{zn} \right) e^{-jnp\theta} \mathbf{i}_r$$

$$\mathbf{B}_{\text{tz}}^{Z_a} = \left\{ \mathbf{B}_z^{\text{II}}(r, \theta, Z_a) + \mathbf{B}_z^{\text{I}}(r', \theta', 0) \right\}$$

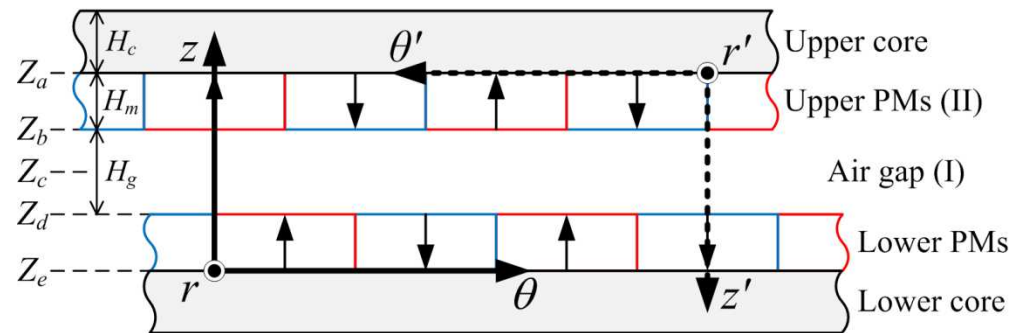
$$\mathbf{B}_z = \sum_{n=-\infty, \text{odd}}^{\infty} \frac{jnp}{r} A_n(z) e^{-jnp\theta} \mathbf{i}_z$$

$$\mathbf{B}_{\text{t}\theta}^{Z_a} = \left\{ \mathbf{B}_\theta^{\text{II}}(r, \theta, Z_a) + \mathbf{B}_\theta^{\text{I}}(r', \theta', 0) \right\}$$

$$\mathbf{B}_\theta = \sum_{n=-\infty, \text{odd}}^{\infty} \frac{\partial}{\partial z} A_n(z) e^{-jnp\theta} \mathbf{i}_\theta$$

II. Analytical Approach – AFPM Coupling

C. Boundary Conditions



$$\mathbf{B}_0^{\text{II}}(r, \theta, Z_a) = 0$$

$$\mathbf{B}_0^{\text{I}}(r, \theta, Z_e) = 0$$

$$\mathbf{B}_z^{\text{I}}(r, \theta, Z_b) = \mathbf{B}_z^{\text{II}}(r, \theta, Z_b)$$

$$\mathbf{B}_0^{\text{I}}(r, \theta, Z_b) = \mathbf{B}_0^{\text{II}}(r, \theta, Z_b)$$

II. Analytical Approach – AFPM Coupling

D. Results

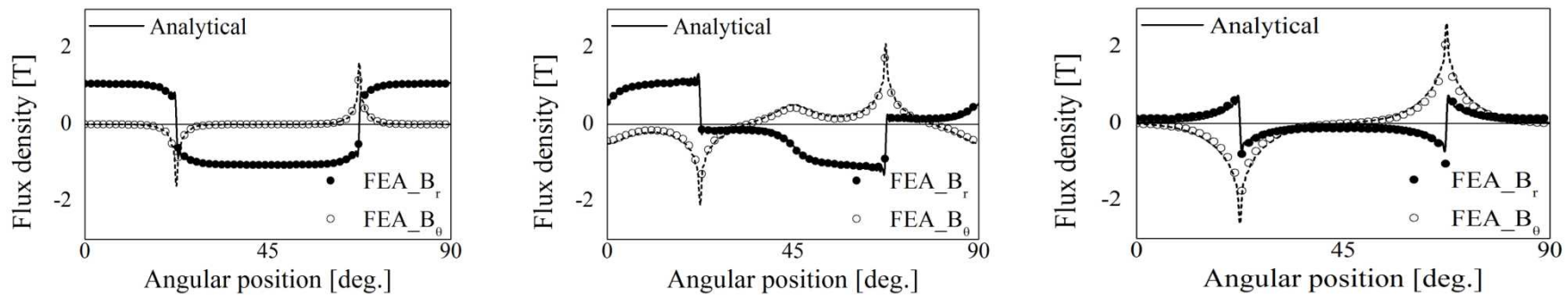


Fig. Comparison of analytical results with 3D FE results for air gap flux density at $g = 4$ mm when θ_a is $0, 22.5^\circ, 45^\circ$.

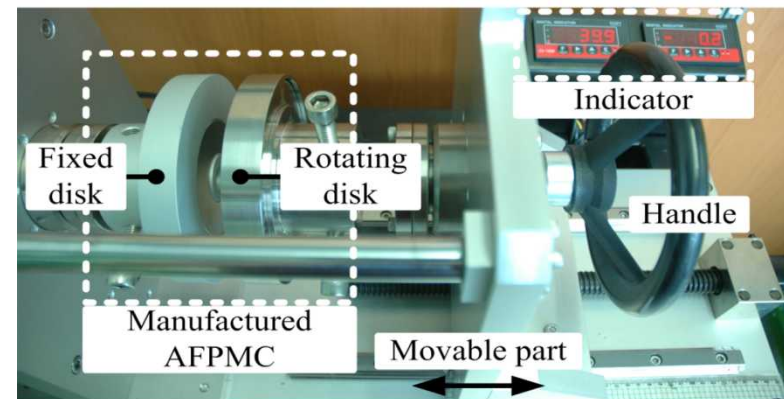


Fig. Testing apparatus for torque measurement of AFPMC.

D. Results

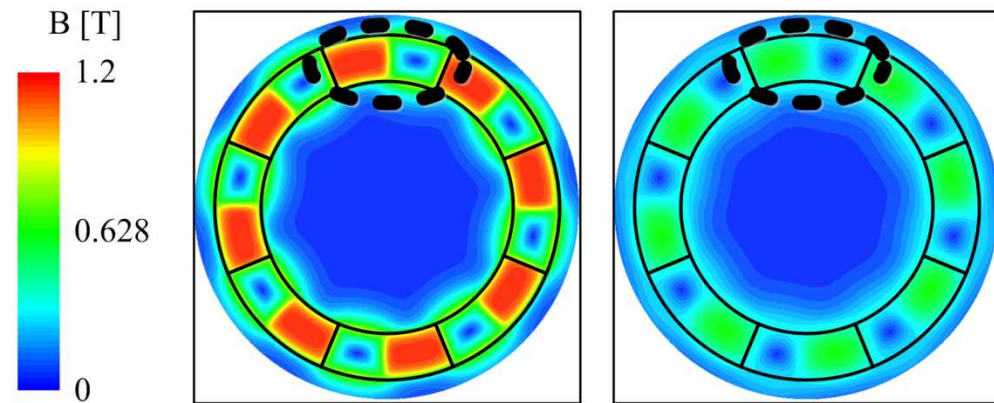
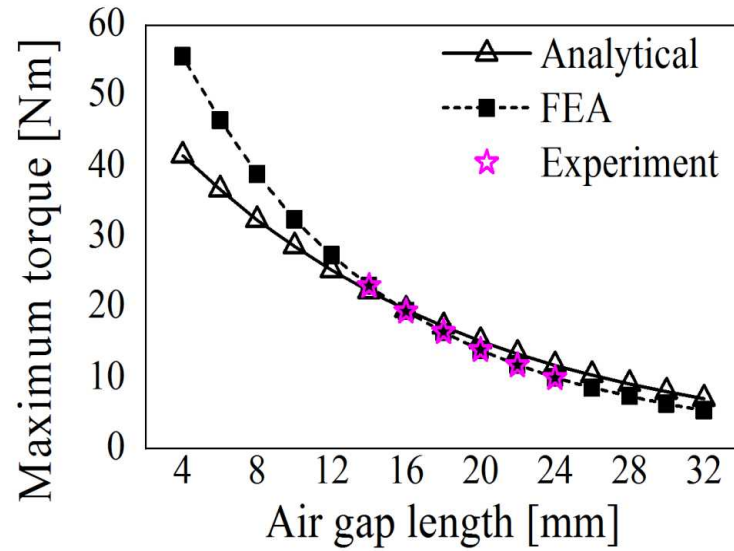


Fig. Magneto-static field distributions obtained from 3D FE analysis at $\theta_a = 22.5^\circ$; (a) $g = 4$ mm and (b) $g = 16$ mm.

III Analytical Approach - Gears

III. Analytical Approach – Gears

A. Analytical Model

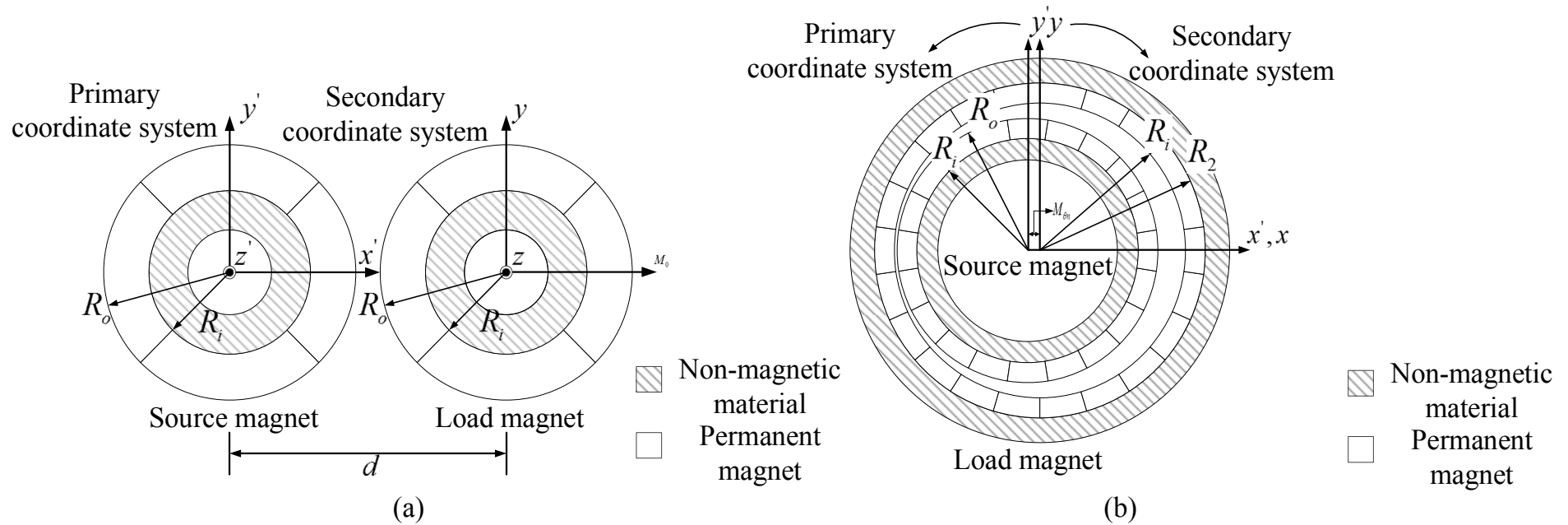


Fig. 6 Schematic diagram of (a) magnetic spur gear (b) magnetic cycloid gear

- The analysis of magnetic gear is implemented in three parts :
 - ① Analysis of magnetic fields produced by source magnet
 - ② Coordinate conversion
 - ③ Torque calculation by using Lorentz force

B. Analytical model

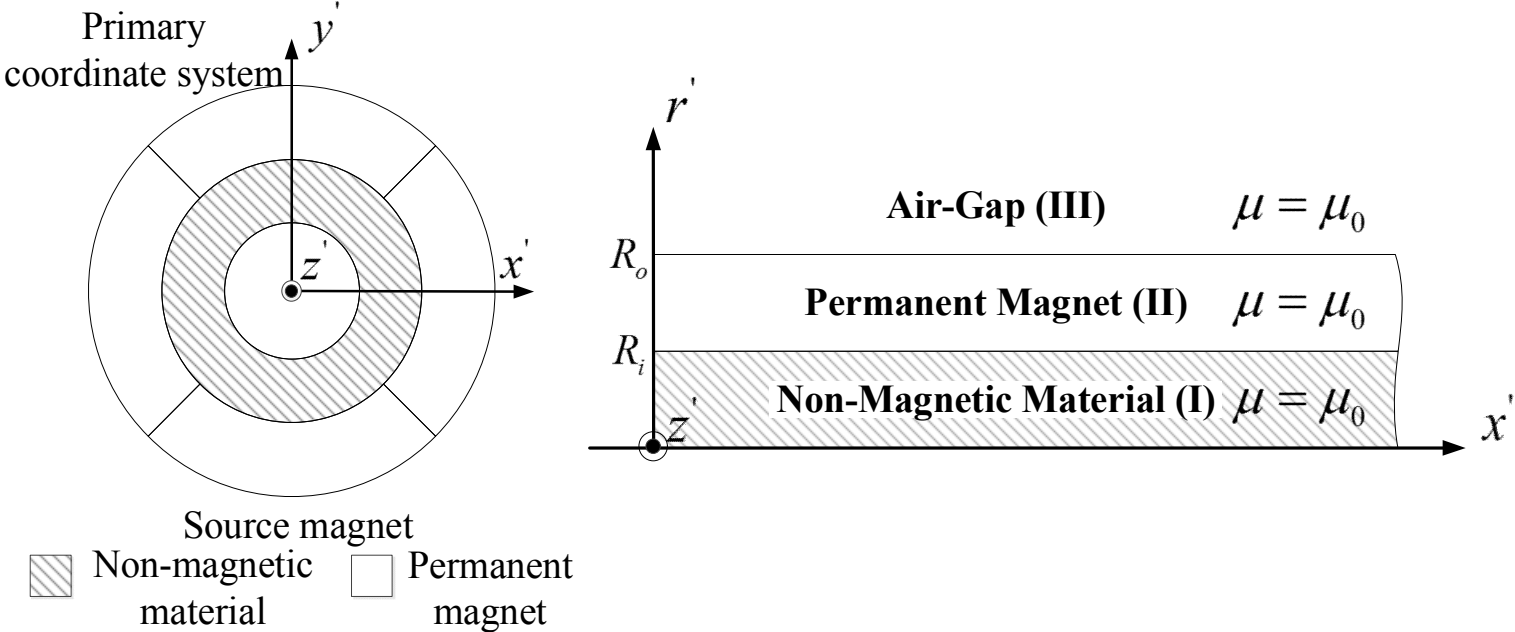


Fig. Magnetic field analysis model of the source magnet

C. Magnetization model

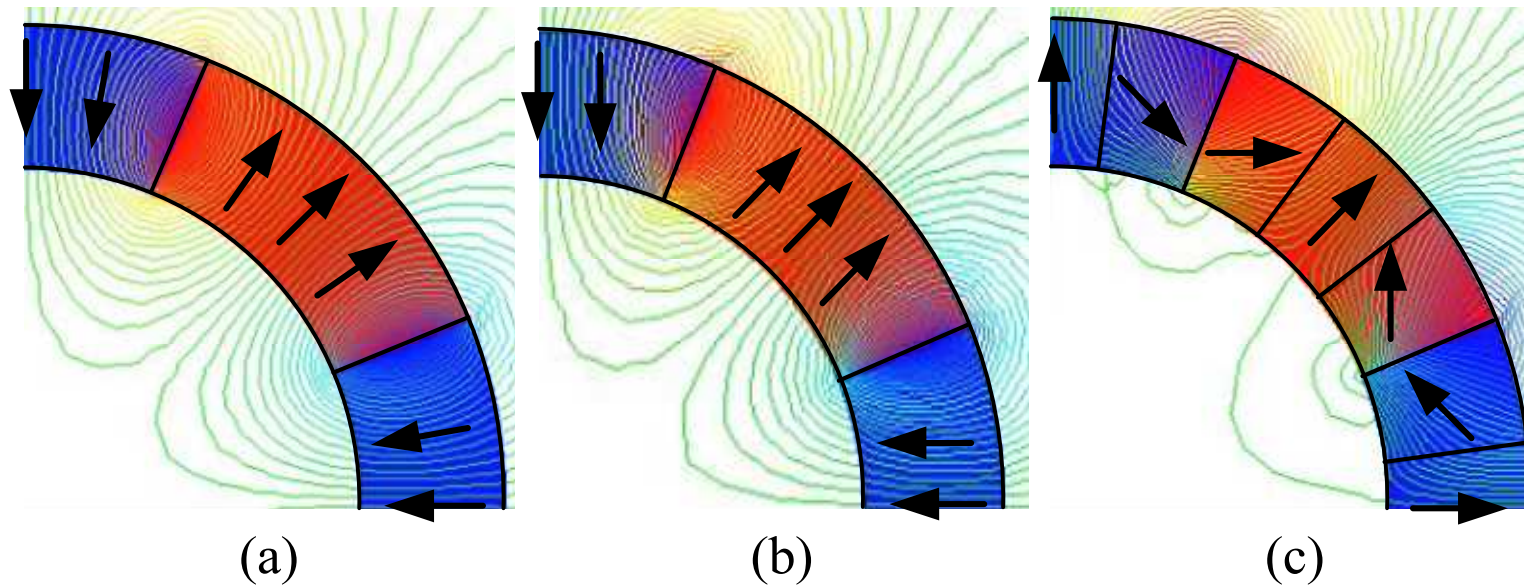


Fig. Magnetization of magnets : (a) radial magnetization (b) parallel magnetization (c) 3-segments Halbach magnetization

III. Analytical Approach – Gears

D. Magnetic field calculations – governing equations

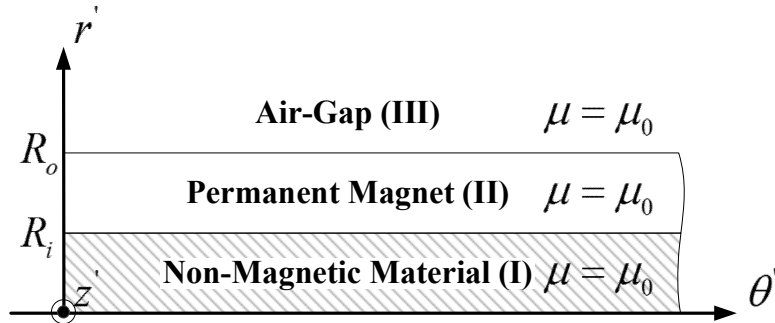


Fig. 12 Analytical model

$$\mathbf{B} = \mu_0 (\mathbf{H} + \mathbf{M})$$

$$\nabla \times \mathbf{B} = \nabla \times \mu_0 \mathbf{M},$$

$$\mathbf{B} \equiv \nabla \times \mathbf{A}$$

$$\nabla \cdot \mathbf{A} = 0$$

$$\nabla^2 \mathbf{A}^{I,III} = 0$$

$$\nabla^2 \mathbf{A}^{II} = -\mu_0 (\nabla \times \mathbf{M})$$

$$\mathbf{A} = \sum_{n=-\infty, \text{odd}}^{\infty} A_n(r) e^{-jnp_s(\theta')} \cdot \mathbf{i}_z$$

$$\mathbf{A}^{I,III} = \sum_{n=-\infty, \text{odd}}^{\infty} \left\{ C_n^{I,III} r^{np_s} + D_n^{I,III} r^{-np_s} \right\} \cdot e^{-jnp_s \theta'} \mathbf{i}_z$$

$$\mathbf{A}^{II} = \sum_{n=-\infty, \text{odd}}^{\infty} \left\{ \begin{array}{l} C_n^{II} r^{np_s} + D_n^{II} r^{-np_s} \\ -\frac{\mu_0 (jnp_s M_{rn} + M_{\theta n})}{1 - (np_s)^2} r \end{array} \right\} \cdot e^{-jnp_s \theta'} \mathbf{i}_z$$

$$\mathbf{B} \equiv \nabla \times \mathbf{A}$$

$$\mathbf{B}_r^{III} = -\sum_{n=-\infty}^{\infty} \frac{jnp_s}{r'} A_n^{III}(r') e^{-jnp_s \theta'} \mathbf{i}_r$$

$$\mathbf{B}_{\theta'}^{III} = -\frac{\partial}{\partial r} \sum_{n=-\infty}^{\infty} \frac{jnp_s}{r'} A_n^{III}(r') e^{-jnp_s \theta'} \mathbf{i}_{\theta}$$

D. Magnetic field calculations – boundary conditions

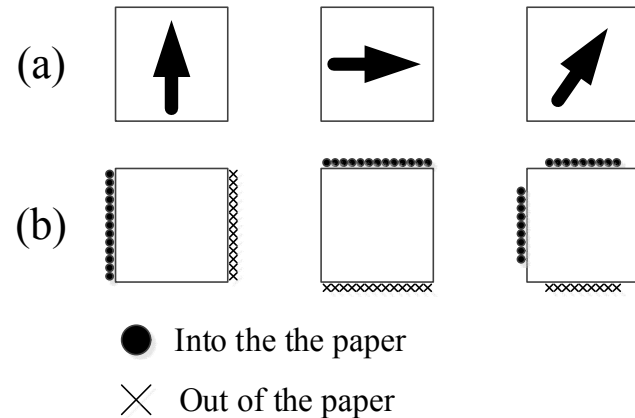


Fig. Electromagnetic dual of any magnetization : (a) direction of magnetization (b) equivalent current model

$$\begin{aligned}
 & 1 \quad \lim_{r \rightarrow 0} A_n^I(r, \theta) = 0 \quad \text{and} \quad \lim_{r \rightarrow \infty} A_n^{III}(r, \theta) = 0 \\
 &) \\
 & 2 \quad B_{rn}^I(R_i, \theta') = B_{rn}^{II}(R_i, \theta') \quad \text{at } r = R_i \\
 &) \quad B_{\theta n}^I(R_i, \theta') - B_{rn}^{II}(R_i, \theta') = -\mu_o M_{\theta n} \\
 & 3 \quad B_{rn}^{II}(R_o, \theta') = B_{rn}^{III}(R_o, \theta') \quad \text{at } r = R_o \\
 &) \quad B_{\theta n}^{II}(R_o, \theta') - B_{rn}^{III}(R_o, \theta') = \mu_o M_{\theta n}
 \end{aligned}$$

J. Y. Choi, H. Y. Kim, S. M. Jang, S. H. Lee, "Thrust Calculations and Measurements of Cylindrical Linear Actuator Using Transfer Relations Theorem," *IEEE Trans. Magn.*, vol.44, pp.4081-4084, Nov. 2008.

E. Coordinate conversion

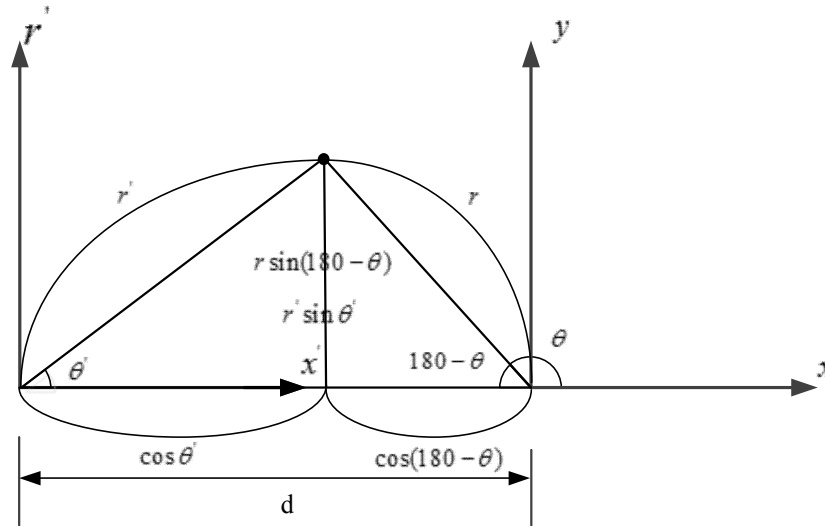


Fig. Schematic of coordinate conversion

$$\begin{aligned}
 r' \sin \theta' &= r \sin(180 - \theta) \\
 d &= r' \cos \theta' + r \cos(180 - \theta)
 \end{aligned}
 \quad \longrightarrow \quad
 \begin{aligned}
 r' &= \sqrt{r^2 + 2rd \cos(\theta) + d^2} \\
 \theta' &= \arctan \left(\frac{r \sin(\theta)}{r \cos(\theta) + d} \right)
 \end{aligned}$$

E. Coordinate conversion

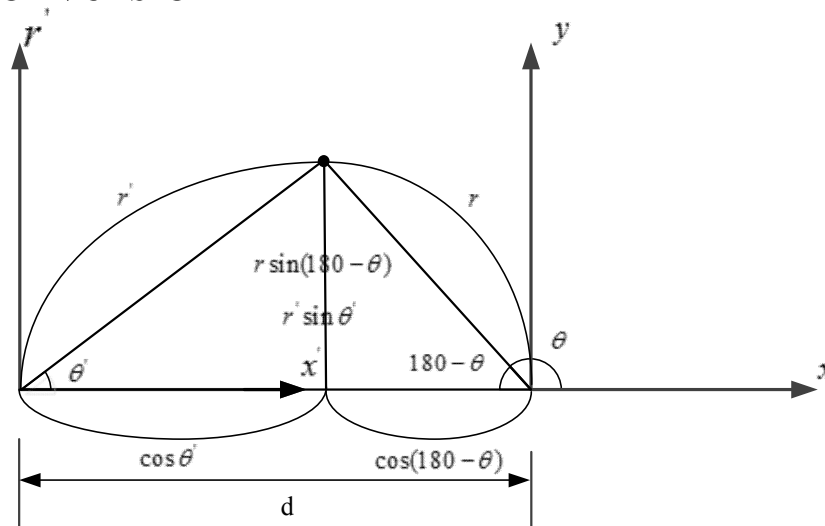


Fig. Schematic of coordinate conversion

$$\begin{matrix}
 B_r^{III}(r', \theta') \\
 B_\theta^{III}(r', \theta')
 \end{matrix}
 \rightarrow
 \begin{matrix}
 B_r^{III}(r, \theta) \\
 B_\theta^{III}(r, \theta)
 \end{matrix}$$

$$\begin{aligned}
 B_x^{ext} &= B_r^{III}(r, \theta) \cdot \cos \left\{ \arctan \left(\frac{r \sin \theta}{r \cos \theta + d} \right) \right\} \\
 &\quad - B_\theta^{III}(r, \theta) \sin \left\{ \arctan \left(\frac{r \sin \theta}{r \cos \theta + d} \right) \right\} \\
 B_y^{ext} &= B_r^{III}(r, \theta) \cdot \sin \left\{ \arctan \left(\frac{r \sin \theta}{r \cos \theta + d} \right) \right\} \\
 &\quad + B_\theta^{III}(r, \theta) \cos \left\{ \arctan \left(\frac{r \sin \theta}{r \cos \theta + d} \right) \right\}
 \end{aligned}$$

The Cartesian coordinate is common to both systems

F. Torque Analysis

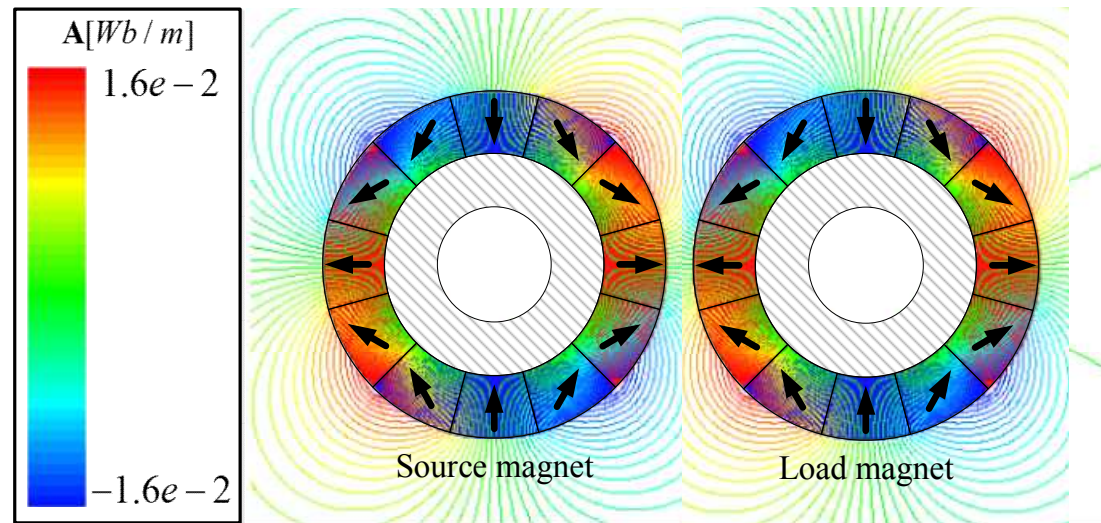


Fig. Torque computed on currents in the external field

$$d\mathbf{T} = \mathbf{r} \times (\mathbf{J} \times \mathbf{B}_{ext}) dV$$

$$\mathbf{T} = \int_S \mathbf{r} \times (\mathbf{j}_m \times \mathbf{B}_{ext}) da$$

$$\mathbf{j}_m = \mathbf{M} \times \mathbf{n}$$

$$\mathbf{J}_m = \nabla \times \mathbf{M} = 0$$

F. Torque Analysis

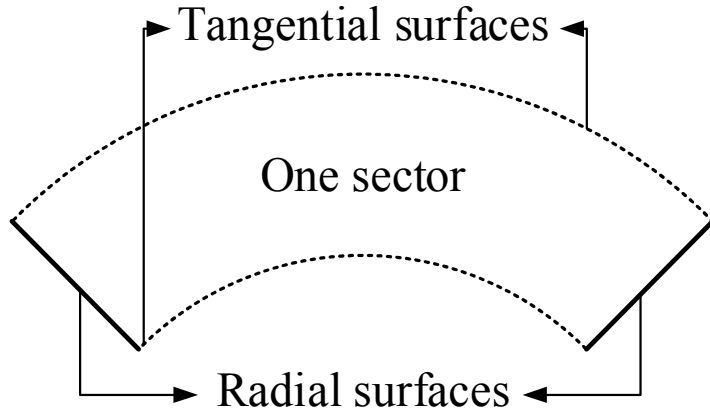


Fig. The radial surface torque and the tangential surface torque

$$T(\theta) = T_r(\theta) + T_{t1}(\theta) + T_{t2}(\theta)$$

$$T_r(\theta) = 2M \cos\left(\frac{\pi}{2p_l}\right) L \frac{(R_2 - R_1)}{N_r} \sum_{p=0}^{2p_l-1} \sum_{q=0}^{N_r} (-1)^p S_r(q) r(q) \times [\cos(\theta_{edge}(\theta, p)) B_x^{ext}(r, \theta_{edge}(\theta, p)) + \sin(\theta_{edge}(\theta, p)) B_y^{ext}(r, \theta_{edge}(\theta, p))]$$

$$T_{t1}(\theta) = -ML \frac{R_1\left(\frac{\pi}{p_l}\right)}{N_t} \sum_{p=0}^{2p_l-1} \sum_{q=0}^{N_t} (-1)^p S_r(q) R_1 \sin(\theta(q)) \times [\cos(\theta(q) + p\frac{\pi}{p_l} + \theta) B_x^{ext}(R_1, \theta(q) + p\frac{\pi}{p_l} + \theta) + \sin(\theta(q) + p\frac{\pi}{p_l} + \theta) B_y^{ext}(R_1, \theta(q) + p\frac{\pi}{p_l} + \theta)]$$

$$T_{t2}(\theta) = ML \frac{R_2\left(\frac{\pi}{p_l}\right)}{N_t} \sum_{p=0}^{2p_l-1} \sum_{q=0}^{N_t} (-1)^p S_r(q) R_2 \sin(\theta(q)) \times [\cos(\theta(q) + p\frac{\pi}{p_l} + \theta) B_x^{ext}(R_2, \theta(q) + p\frac{\pi}{p_l} + \theta) + \sin(\theta(q) + p\frac{\pi}{p_l} + \theta) B_y^{ext}(R_2, \theta(q) + p\frac{\pi}{p_l} + \theta)]$$

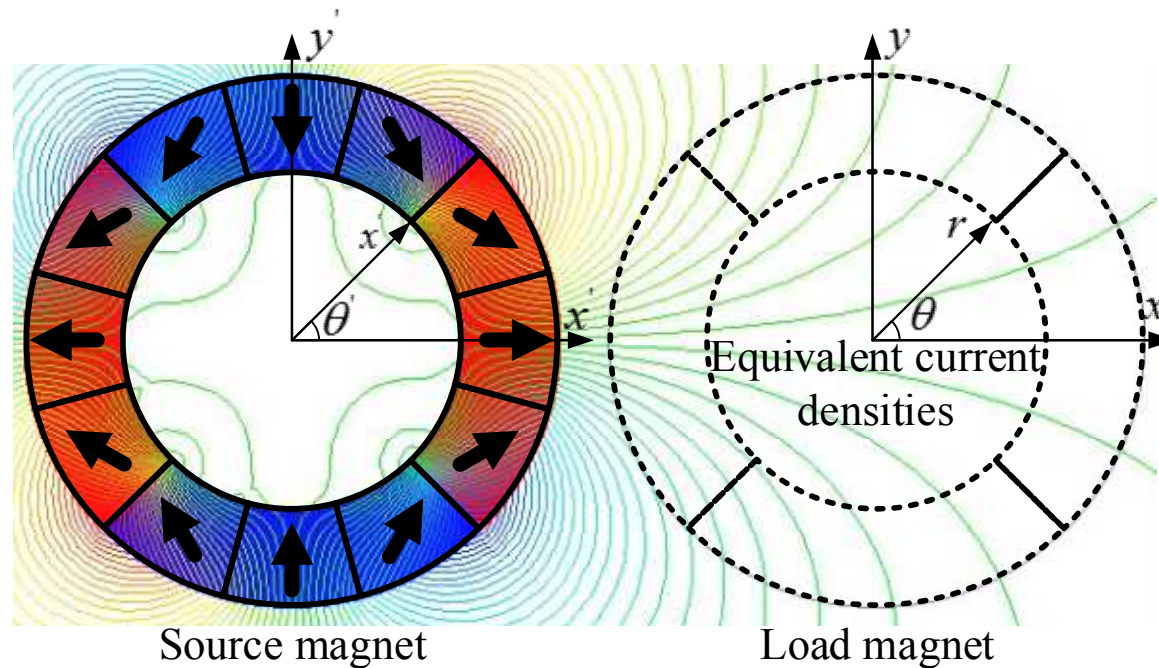


Fig. 21 Magnetic flux line distribution of 3-segments Halbach magnetic spur gear

III. Analytical Approach – Gears

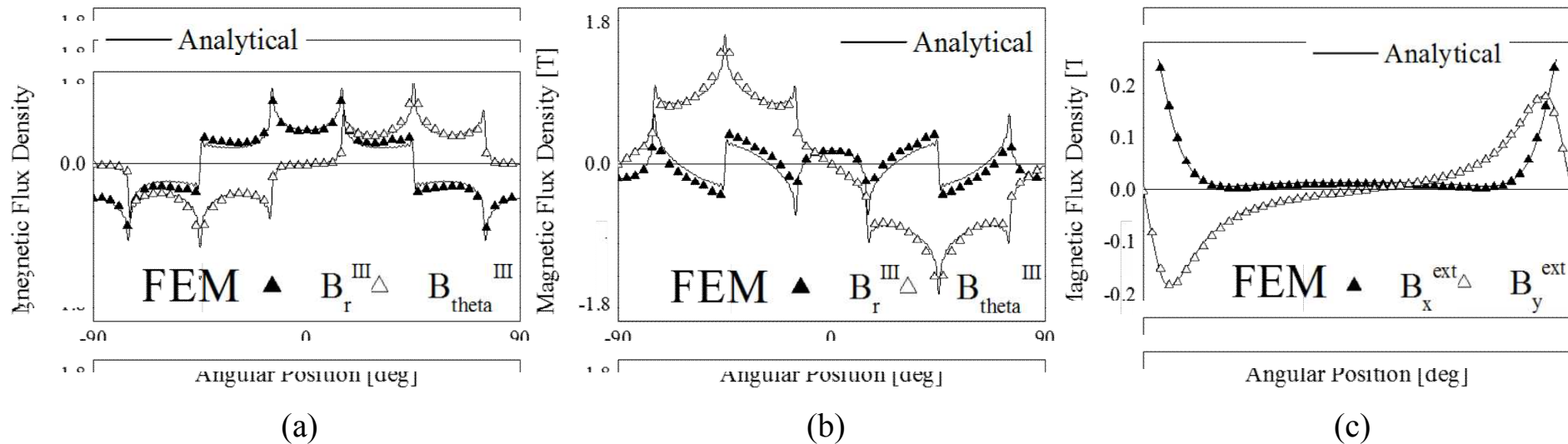


Fig. Comparison of magnetic flux density between the analytical calculations and 2D FE analysis for Halbach spur gear (a) at $r' = 35mm$ (b) at $r' = 55mm$ (c) at $r = 60mm$

III. Analytical Approach – Gears

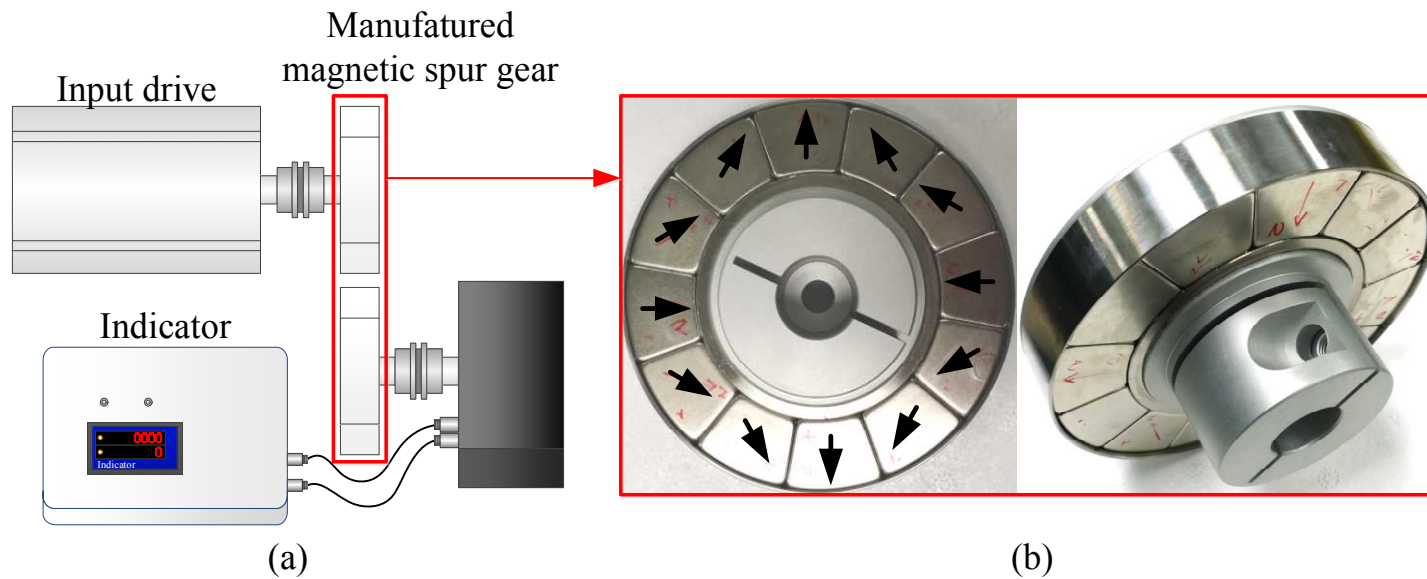


Fig. (a) Test apparatus for torque measurement (b) Manufactured magnetic spur gear with Halbach magnetized PMs.

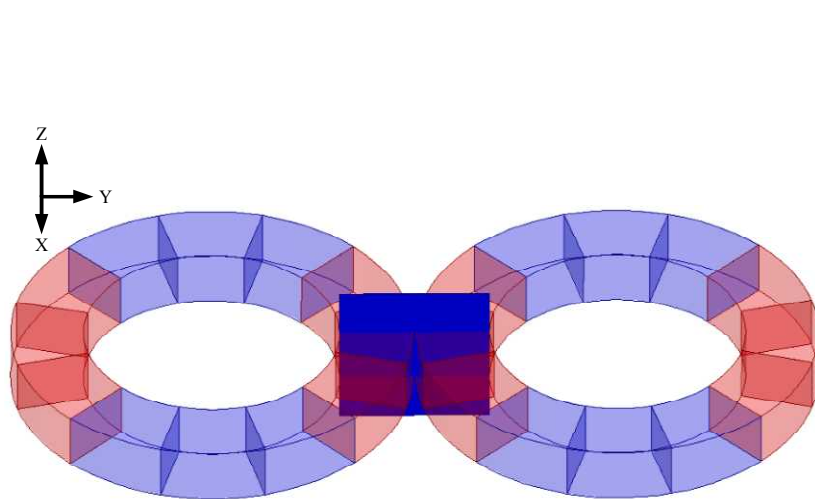


Fig. 3D FE analysis model of magnetic spur gear

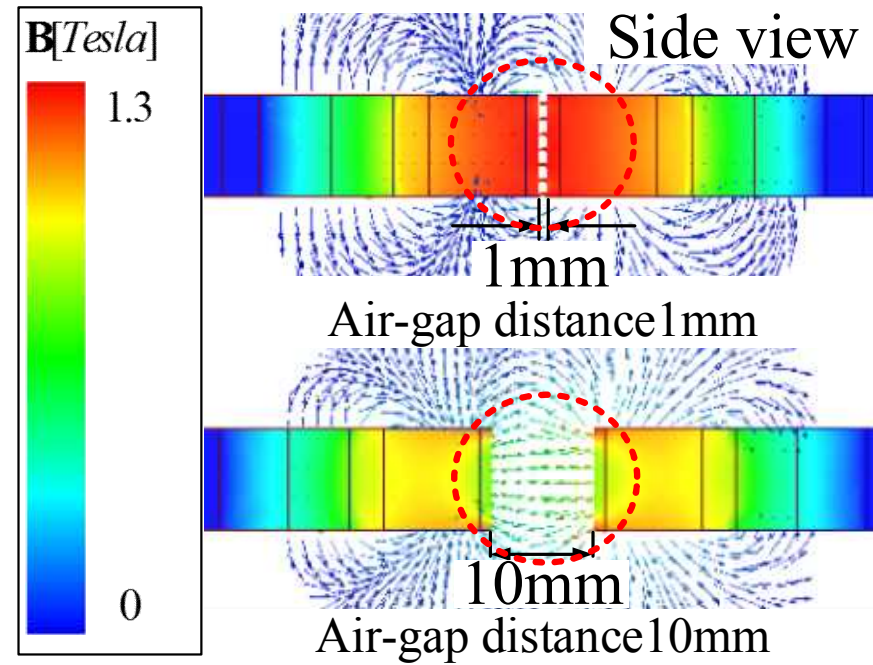


Fig. 3D effect on magnetic spur gear with 3-segments Halbach magnetized PMs

Magnetization	Air-gap distance	Analytical result (error)	2D FE result (error)	Measurement
Halbach	1mm	13 [Nm] (26%)	13.4 [Nm] (30%)	10.29 [Nm]
Halbach	10mm	9.4 [Nm] (52%)	9.9 [Nm] (60%)	6.17 [Nm]

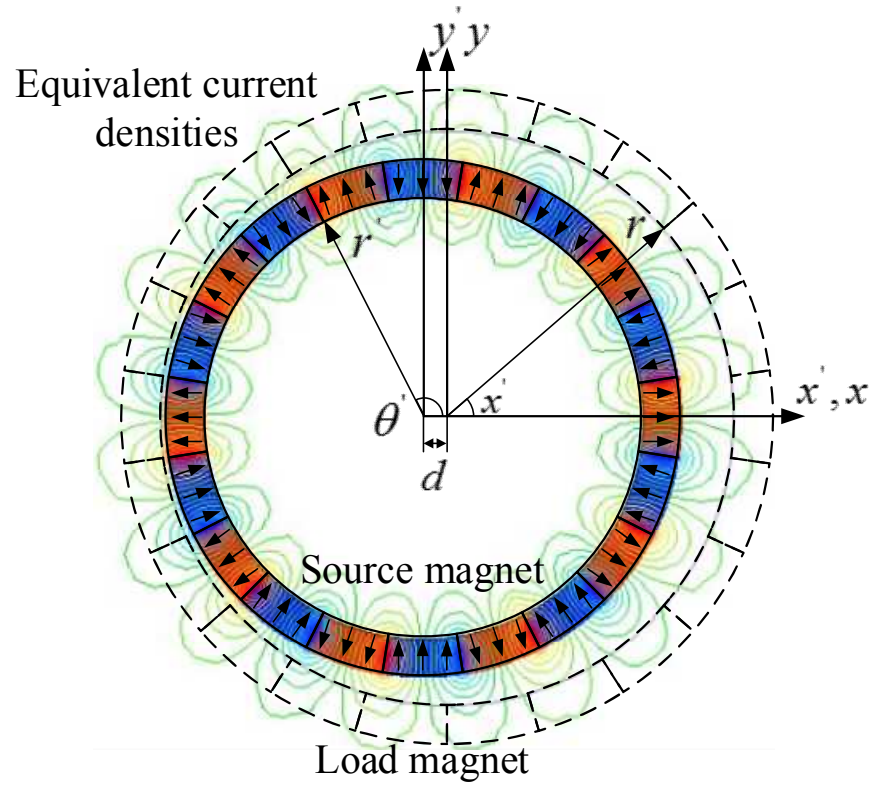
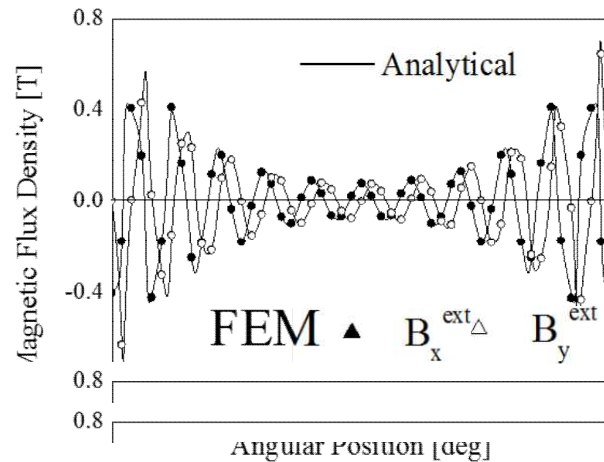
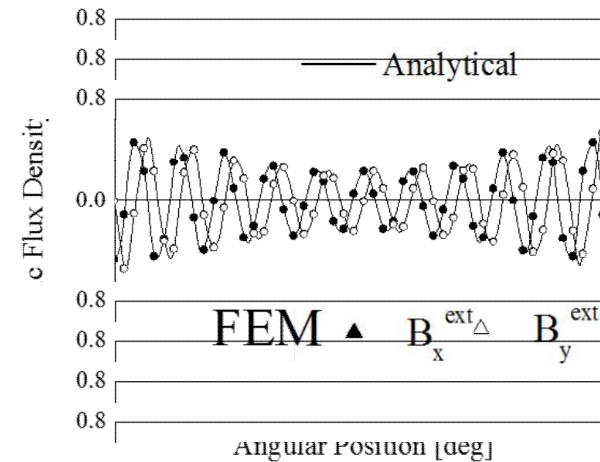


Fig. 28 Magnetic flux line distribution of parallel magnetic cycloid gear $d = 2.5mm$

III. Analytical Approach – Gears



(a)



(b)

Fig. Comparison of magnetic flux density of magnetic cycloid gear with parallel magnetized PMs between the analytical calculations and 2D FEA when $r = 30mm$, (a) $d = 2.5mm$ (b) $d = 1mm$

III. Analytical Approach – Gears

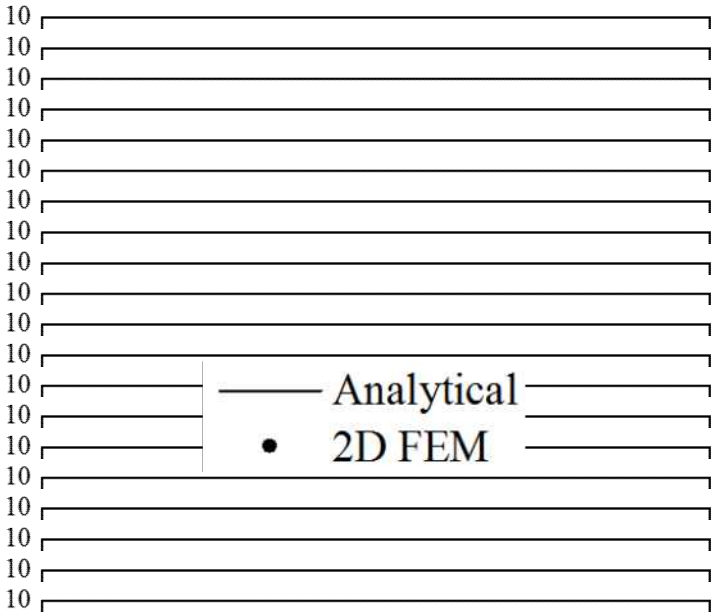


Fig. Comparison of the torque calculation and FE results, when $d = 2.5mm$

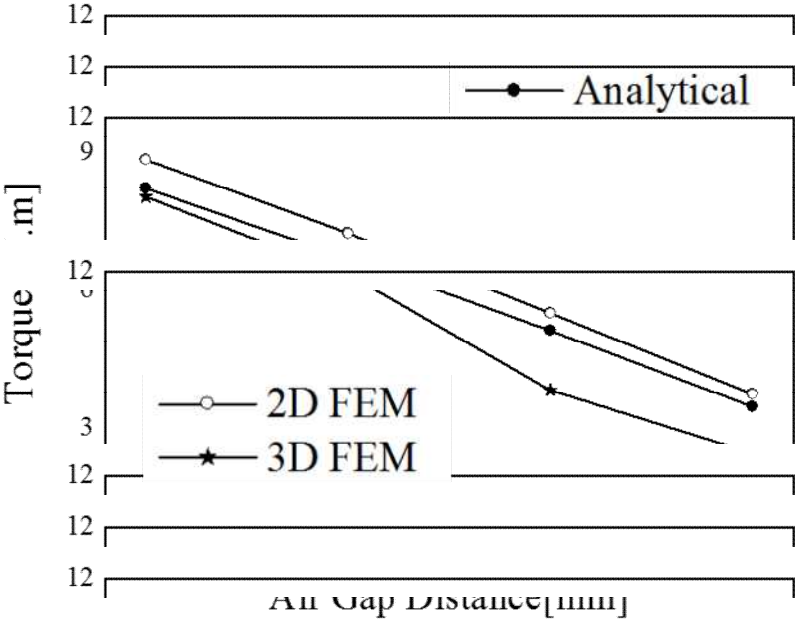


Fig. Comparison of the torque calculation and FE results, when $1mm \leq d \leq 2.5mm$

IV Conclusion

- The analytical torque calculation for a magnetic and coupling is presented and compared with the 2D, 3D FE and experimental results
- The analytical results are shown in good agreement with the 2D FEA results.
- However, the analytical results are different from experimental results and 3D FE results
- Difference of torque results are caused by 3D effects
- The percentage of error between the analytical results and measurement, 3D FE results should increase when the air-gap distance is enlarged.
- In spite of the 3D effects, It would be possible to use these analytical calculations for initial design and optimization purposes.

Thank you
for your attention